

Adiabatic shear sensitivity of ductile metal based on gradient-dependent JOHNSON-COOK model

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Abstract: Based on the expression proposed by WANG for the local plastic shear deformation distribution in the adiabatic shear band (ASB) using gradient-dependent plasticity, the effects of 10 parameters on the adiabatic shear sensitivity were studied. The experimental data for a flow line in the ASB obtained by LIAO and DUFFY were fitted by use of the curve-fitting least squares method and the proposed expression. The critical plastic shear strains corresponding to the onset of the ASB for Ti-6Al-4V were assessed at different assigned ASB widths. It is found that the proposed expression describes well the non-linear deformation characteristics of the flow line in the ASB. Some parameters in the JOHNSON-COOK model are back-calculated using different critical plastic shear strains. The adiabatic shear sensitivity decreases as initial static yield stress, work to heat conversion factor and strain-rate parameter decrease, which is opposite to the effects of density, heat capacity, ambient temperature and strain-hardening exponent. The present model can predict the ASB width evolution process. The predicted ASB width decreases with straining until a stable value is reached. The famous model proposed by DODD and BAI only can predict a final stable value.

Key words: adiabatic shear band; Ti-6Al-4V; JOHNSON-COOK model; width; gradient-dependent plasticity theory; local plastic shear deformation distribution

1 Introduction

Because titanium alloys possess the properties of high strength, low density and low thermal conductivity, they are very sensitive to adiabatic shear under dynamic loading conditions [1–4]. Adiabatic shear bands (ASBs) are narrow bands in which deformational lines (or called flow lines) are curved (Fig. 1(a)), showing highly nonuniform deformation characteristics. Curved flow lines are observed in the later stage of the plateau stage of the shear stress-nominal shear strain curve (Fig. 1(b)). The flow lines in Fig. 1(a) correspond to the position of number 3 in Fig. 1(b). Using traditional elastoplastic theories, the distribution of the strain, deformation, temperature and damage extent in the ASB and the band width cannot be predicted.

The ASB width can provide a measure of the adiabatic shear sensitivity. Some experimental results show that ASBs become narrow at high strain rates and low ambient temperatures. To model the thickness of the ASB for Ti-6Al-4V at different loading rates [5–7], WANG [8] proposed a simple model. In the model, the

ASB width was defined as the width ($w_{5\%}$) of the region surrounding the band center over which the temperature differed from its peak value by less than 5%, which was similar to the viewpoint proposed by BATRA and KIM [9]. In the proposed model, the temperature distribution in the ASB was nonuniform [10–11], which was dependent on the plastic shear strain distribution [12–13]. The plastic shear strain distribution and the total thickness of the ASB were controlled by the internal length parameter in gradient-dependent plasticity (GDP) [14–15]. Based on the gradient-dependent JOHNSON-COOK model, the width $w_{5\%}$ was calculated [8]. It is shown that the ASB becomes wide at low strain rates; the width $w_{5\%}$ is decreased rapidly as strain rate increases; and the width $w_{5\%}$ remains a constant at high strain rates. The theoretical result is very consistent with the experimental result for tungsten alloy by WEERASOORIYA and BEAULIEU [7].

The present work based on the WANG's model is to study the effects of parameters on the ASB width $w_{5\%}$. The experimental data of the local shear deformation distribution in an ASB for Ti-6Al-4V [3] are fitted by use of the curve-fitting least squares method. Using the

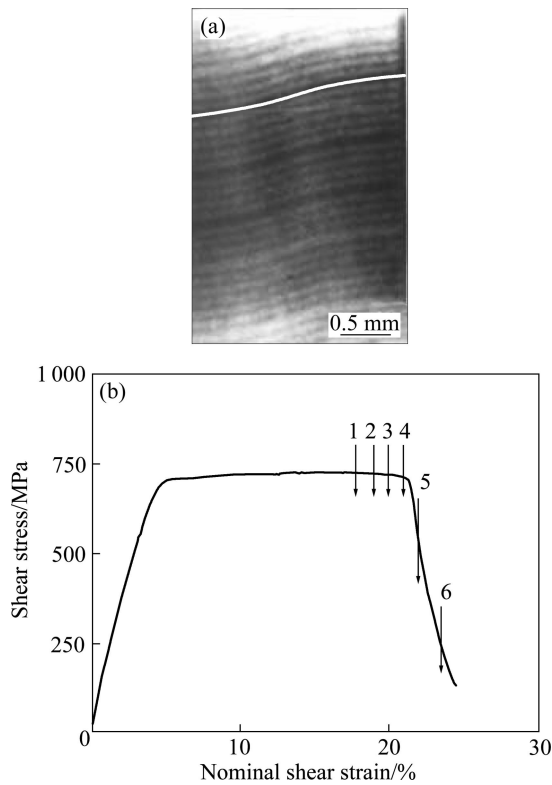


Fig. 1 Flow lines for Ti-6Al-4V specimen using high-speed photography (a) and measured shear stress—shear strain curve (b) in dynamic torsion test [3]

different critical plastic shear strains, some parameters in the JOHNSON-COOK model are back-calculated and the influence of some parameters on the adiabatic shear sensitivity is investigated.

2 Theoretical analyses

2.1 Local plastic shear deformation distribution in ASB and coordinate transformation

A narrow ASB subjected to shear stress is shown in Fig. 2. A simple expression for the local plastic shear deformation distribution $\gamma_p(y)$ in an ASB was proposed by WANG [8, 10–12] using GDP:

$$\gamma_p(y) = \gamma_c^p + (\bar{\gamma}_p - \gamma_c^p) \cdot \left(1 + \cos \frac{y}{l}\right) \quad (1)$$

where y is the coordinate whose original point is set at the ASB center; γ_c^p is the critical plastic shear strain corresponding to the onset of the ASB; if the ASB is assumed to be initiated just at the peak shear stress, then γ_c^p is equal to the average plastic shear strain corresponding to the peak shear stress; $\bar{\gamma}_p$ is the average plastic shear strain of the ASB, which is a variable and is increased with loading time; and l is the internal length parameter in GDP, reflecting the heterogeneous extent of ductile metal material. For

material with a higher internal length parameter, it is more nonuniform in microstructures. The internal length parameter depends on the total thickness w of the ASB: $l = w/(2\pi)$.

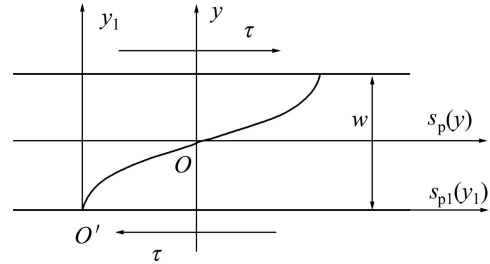


Fig. 2 ASB subjected to shear stress

Integrating $\gamma_p(y)$ leads to the local plastic shear deformation distribution $s_p(y)$ in the ASB:

$$s_p(y) = \bar{\gamma}_p y + l(\bar{\gamma}_p - \gamma_c^p) \sin \frac{y}{l} \quad (2)$$

It can be found that the expression for $s_p(y)$ includes four parameters: $\bar{\gamma}_p$, l , γ_c^p and y . For a flow line in high-speed photographs, $\bar{\gamma}_p$, l and γ_c^p are specific, while y is a variable denoting different positions in the ASB. Introducing two coefficients c_1 and c_2 leads to:

$$s_p(y) = c_1 y + c_2 \sin \frac{y}{l} \quad (3)$$

where $c_1 = \bar{\gamma}_p$ and $c_2 = l(\bar{\gamma}_p - \gamma_c^p)$. The expression for γ_c^p is expressed by c_1 and c_2 :

$$\gamma_c^p = c_1 - \frac{c_2}{l} \quad (4)$$

In Eq. (2), the original point is set at the ASB center O . For a flow line in the ASB, the determination of the ASB center is not easy. Therefore, a coordinate transformation is necessary. The coordinate axis y is moved to axis y_1 and the coordinate axis $s_p(y)$ is moved to $s_{p1}(y_1)$. Thus, the new original point is O' (Fig. 2). In the point O' , the local plastic shear deformation $s_{p1}(0)$ is zero and local plastic shear deformations of other points are the relative values to the point O' .

Apparently, the coordinates y , y_1 , $s_p(y)$ and $s_{p1}(y_1)$ have the following relations:

$$y_1 = y + \frac{w}{2} \quad (5)$$

$$s_{p1}(y_1) = s_p(y) + c_1 \frac{w}{2} \quad (6)$$

After substituting Eqs. (3) and (5) into Eq. (6), the expression for $s_{p1}(y_1)$ is

$$s_{p1}(y_1) = c_1 y_1 + c_2 \sin \frac{y_1 - \pi l}{l} \quad (7)$$

2.2 Curve-fitting of experimental data based on theoretical expression

The two coefficients c_1 and c_2 in Eq. (7) can be obtained by the curve-fitting least squares method. The value of the length parameter l needs to be determined by the total thickness w of the ASB. Obtaining optimal c_1 and c_2 is the aim of the curve-fitting least squares method. They make the sum of the difference squared between the experimental data (y_{1i}, s_i) and theoretical results minimum:

$$\min \sum_{i=1}^N [s_i - s_{pl}(c_1, c_2, y_{1i})]^2 \quad (8)$$

where $i=1, 2, \dots, N$, and N is the number of data.

2.3 Central region thickness $w_{5\%}$ of ASB

In the ASB, the temperature is nonuniform, as shown in Fig. 3. WANG [8] introduced an assumption in the determination of the central region thickness $w_{5\%}$ of the ASB. It is similar to the viewpoint proposed by BATRA and KIM [9]. Using the GDP and JOHNSON-COOK model, the width $w_{5\%}$ for Ti-6Al-4V was calculated and the strain rate effect was studied. The deduced width $w_{5\%}$ was [8]

$$w_{5\%} = 2l \arccos \left(\frac{\frac{19}{10} \int_{\gamma_c}^{\bar{\gamma}_p} \tau d\bar{\gamma}_p - \frac{T_0 \rho c_p}{20\beta} - \frac{1}{20} \int_0^{\gamma_c} \tau d\bar{\gamma}_p}{\int_{\gamma_c}^{\bar{\gamma}_p} \tau d\bar{\gamma}_p} - 1 \right) \quad (9)$$

where τ is the shear stress acting on the ASB (Fig. 2); T_0 is the ambient temperature or initial temperature (Fig. 3); β is the work to heat conversion factor; c_p is the heat capacity; and ρ is the density. In Fig. 3, T_1 is $\int_0^{\gamma_c} \tau d\bar{\gamma}_p$; $T(y)$ is the temperature distribution in the ASB; and $T(0)$ is peak temperature.”

In the JOHNSON-COOK model, the shear stress τ and temperature T were expressed as [8, 10–11, 13]

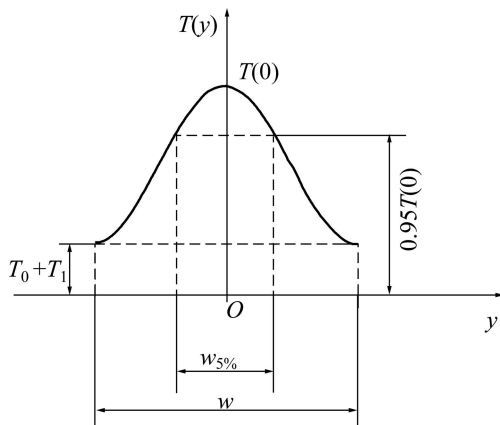


Fig. 3 Temperature distribution in ASB and definition of ASB width $w_{5\%}$

$$\tau = (A + B\bar{\gamma}_p^n) \cdot \left(1 + C \ln \frac{\dot{\gamma}}{\dot{\gamma}_0} \right) \cdot \left[1 - \left(\frac{T - T_0}{T_m - T_0} \right)^m \right] \quad (10)$$

$$T = T_0 + \frac{\beta}{\rho c_p} \int \tau d\bar{\gamma}_p \quad (11)$$

where A is the initial static yield stress; B is the strain-hardening coefficient; C is the strain-rate parameter; m is the thermal-softening parameter; n is the strain-hardening exponent; T_m is the melting temperature; $\dot{\gamma}_0$ is the reference shear strain rate; and $\dot{\gamma}$ is the exerted shear strain rate.

3 Results and discussion

3.1 Curve-fitting of experimental flow line data

The process of initiation and formation of ASBs in Ti-6Al-4V alloy was investigated in dynamic torsional experiments with strain rate of about 10^3 s^{-1} at room temperature [3]. In these tests, thin-wall tubular specimens were deformed dynamically in a torsional Kolsky bar and high-speed photography was used. For the specimen TA-43, six high-speed photographs of the grid pattern were taken during the formation of an ASB [3]. In these photographs, some curved flow lines can be found. These lines are oriented parallel to the axis of specimen before experiment. Some flow lines are bold and clear. The fitted data are taken in a bold flow line in Fig. 1(a). The flow line has been painted white. In the white line, a total of 41 data from left to right is selected. The horizontal distance between two adjacent data points is 0.05 mm.

Firstly, all of 41 data are fitted by use of Eqs. (7)–(8). The first data point is selected as the original point O' . Thus, the total thickness of the ASB is assigned to be 2 mm. Then, a part of data, i.e., 39, 37, 35, 33, 31, 29, 27 and 25, are selected and fitted, respectively. When 39 is selected, the first and last data are omitted. When 37 is selected, the former two and last two data are omitted, and so on.

The total thickness w of the ASB can be calculated by the number N of data:

$$w = 0.05 \times (N - 1) \quad (12)$$

A total of nine times curve-fitting is conducted in the work. Only the six times results are plotted in Fig. 4. Figure 4 shows that the agreement between the experimental data and theoretical results is very good at different N . It is found that the theoretical expression for the local plastic shear deformation distribution in the ASB can well describe the nonlinear characteristics of the flow line. When the total thickness is assigned to be a higher value, the flow line in the ASB is more curved at the two ends of the band, while when the total thickness

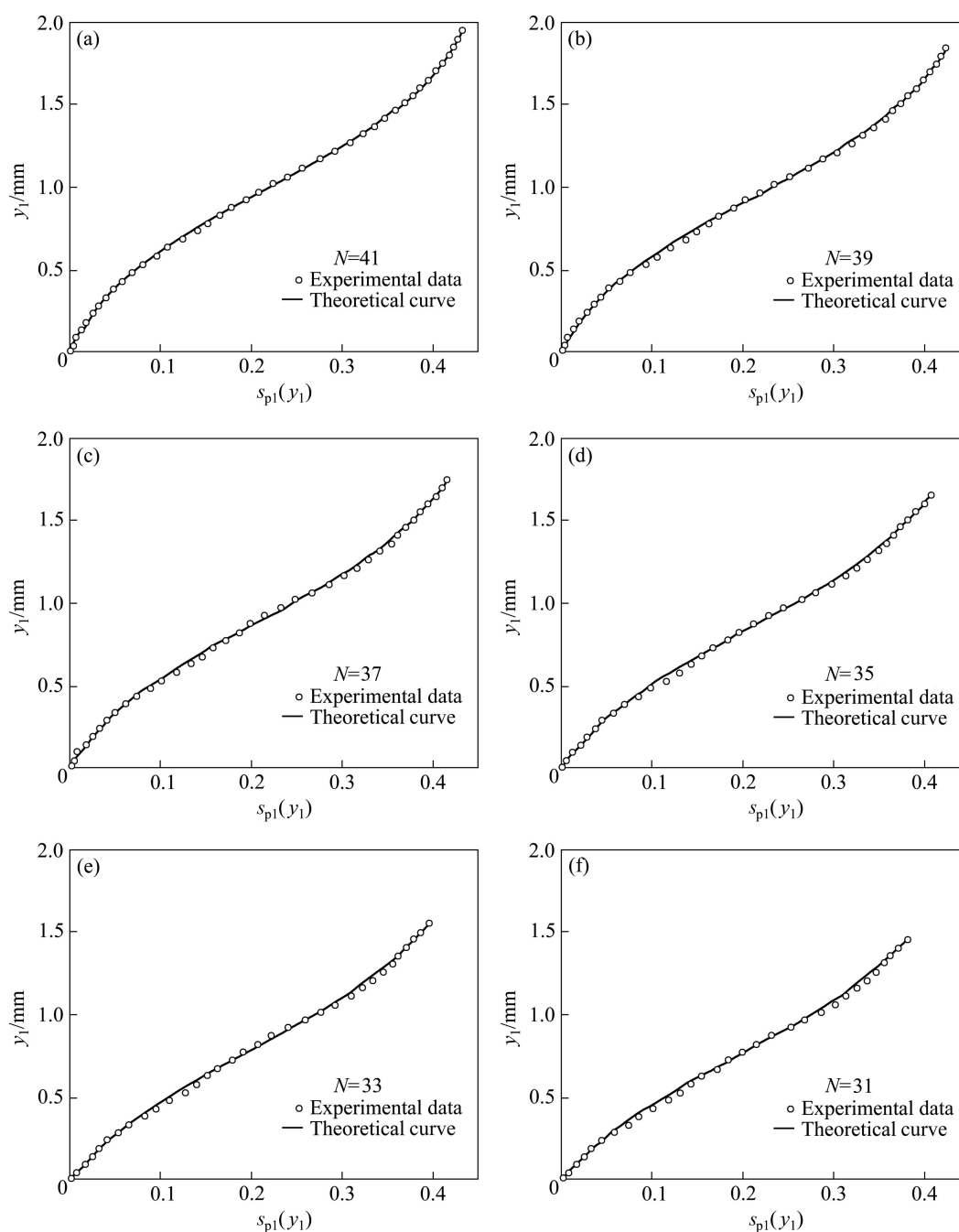


Fig. 4 Comparisons of experimental data [3] and theoretical results for local deformation distribution in ASB at different number

is assigned to be a lower value the flow line is straighter.

Table 1 shows the obtained γ_c^p , $\bar{\gamma}_p$ and residual errors E_r by curve-fitting. It can be found that γ_c^p and E_r become high as N is low.

3.2 Back-calculation of parameters in JOHNSON-COOK

The JOHNSON-COOK model needs some parameters, such as A , B , C , m and n . Different sets of parameters predict different γ_c^p and the maximum shear stress (τ_{\max}). Possibly, several sets of parameters predict the same γ_c^p and τ_{\max} . Herein, m , $\dot{\gamma}_0$, $\dot{\gamma}$, T_0 , T_m ,

ρ , β and c_p are assumed to be constants, i.e., $m=0.5$, $T_m=1$ 932 K, $\rho=4$ 430 kg/m³, $T_0=300$ K, $\dot{\gamma}_0=10^{-5}$ s⁻¹, $c_p=564$ J/(kg·K), $\dot{\gamma}=1$ 100 s⁻¹ and $\beta=0.9$ [6]. For the given γ_c^p and τ_{\max} , only four parameters (A , B , C and n) will be back-calculated.

The back-calculation process is briefly introduced as follows. The searching aims (i.e., the values of γ_c^p and τ_{\max}) are given and the possible variation range of each parameter (A , B , C and n) is set. A special set of parameters will be selected by which the calculated new values of γ_c^p and τ_{\max} using the JOHNSON-COOK model are the closest to the initially given values. The

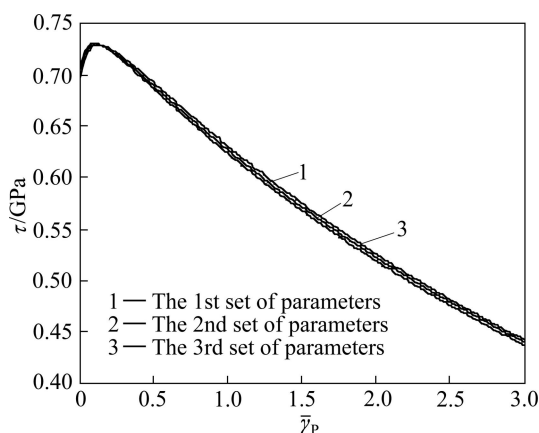
Table 1 Critical plastic shear strains, average plastic shear strains and residual errors obtained by curve-fitting at different number

N	γ_c^p	$\bar{\gamma}_p$	$E_r/10^{-4} \text{ mm}^2$
41	0.104 2	0.219 8	4.870 9
39	0.117 4	0.227 0	4.798 7
37	0.132 9	0.234 9	4.585 4
35	0.151 1	0.243 6	4.267 1
33	0.172 1	0.253 2	4.164 8
31	0.182 5	0.259 2	4.127 1
29	0.195 9	0.266 1	4.058 9
27	0.212 8	0.273 9	3.830 7
25	0.234 4	0.283 0	3.377 8

special set of parameters is the set of back-calculated parameters in the JOHNSON-COOK model.

The value of τ_{\max} has been determined as 729 MPa that corresponds to the nominal shear strain of 12% according to Fig. 1(b). It is found that for the aims of $\gamma_c^p = 0.1$ and $\tau_{\max} = 729$ MPa, the back-calculated parameters are $A = 473.4$ MPa, $B = 482.4$ MPa, $C = 0.003 5$ and $n = 0.198$ (the 1st set of parameters); for the aims of $\gamma_c^p = 0.12$ and $\tau_{\max} = 729$ MPa, the back-calculated parameters are $A = 447.1$ MPa, $B = 482.4$ MPa, $C = 0.005 5$ and $n = 0.2$ (the 2nd set of parameters); and for the aims of $\gamma_c^p = 0.14$ and $\tau_{\max} = 729$ MPa, the back-calculated parameters are $A = 420.8$ MPa, $B = 536$ MPa, $C = 0.004 2$ and $n = 0.19$ (the 3rd set of parameters). The predicted values of τ_{\max} are 730.26 MPa, 729.9 MPa and 728.86 MPa, respectively, which are close to the aim value of 729 MPa using three sets of parameters.

Figure 5 shows the three shear stress—average plastic shear strain curves with three sets of parameters. The results are almost identical. Only minor difference exists. At pre-peak, the highest shear stress is expected using the first set of parameters, while at post-peak this

**Fig. 5** Comparisons of three shear stress—average plastic shear strain curves obtained by three sets of parameters

set of parameters obtains the lowest shear stress.

3.3 Parametric study of ASB width $w_{5\%}$

Herein, only the effects of 10 parameters on the ASB width $w_{5\%}$ are investigated with $N = 41$ (i.e., $w = 2$ mm). The back-calculated parameters ($A = 473.4$ MPa, $B = 482.4$ MPa, $C = 0.003 5$ and $n = 0.198$) when $\gamma_c^p = 0.1$ will be used and the parameters ($m = 0.5$, $T_m = 1 932$ K, $\rho = 4 430$ kg/m³, $T_0 = 300$ K, $\dot{\gamma}_0 = 10^{-5}$ s⁻¹, $c_p = 564$ J/(kg·K), $\dot{\gamma} = 1 100$ s⁻¹ and $\beta = 0.9$) [6] which are seen as constants in the back-calculation process will also be used. When the effect of a parameter is studied, the other parameters are equal to the values given above.

Figure 6 shows the influence of 10 parameters on the ASB width $w_{5\%}$. It can be found that these calculated results can be classified into three categories.

For the first category, the markers with a circle are always above the other two markers at the same average plastic shear strain (Figs. 6(a), (c) and (j)). These results show that the width $w_{5\%}$ increases or adiabatic shear sensitivity decreases as A , β and C decrease.

The second category that is opposite to the first category, shows that the markers with a circle are always below the other two markers at the same average plastic shear strain (Figs. 6(d), (f), (h) and (i)). These results suggest that adiabatic shear sensitivity increases as ρ , c_p , T_0 and n decrease.

For the third category, the markers with a circle are at left side of the other two markers at very low $\bar{\gamma}_p$ (Figs. 6(b), (e) and (g)). When $\bar{\gamma}_p$ reaches a certain value, three kinds of markers are very close. Even, the markers with a circle are above the other two markers at very high $\bar{\gamma}_p$. This suggests that the adiabatic shear sensitivity increases as B , m and T_m decrease at low $\bar{\gamma}_p$; and that at high $\bar{\gamma}_p$ the adiabatic shear sensitivity is nearly independent of the three parameters. Seemingly, the effects of B , m and T_m on the adiabatic shear sensitivity depend on the value of $\bar{\gamma}_p$. However, the effects of A , β , C , ρ , c_p , T_0 and n on the adiabatic shear sensitivity are strictly monotonic.

Figure 6(f) predicts that the adiabatic shear sensitivity increases as the ambient temperature decreases, which is consistent with the experimental results [1, 4]. Figure 6 shows that $w_{5\%}$ tends to be a constant value as $\bar{\gamma}_p$ exceeds a certain value. The reason for the stable ASB width $w_{5\%}$ is possibly due to the stable temperature distribution profile in the ASB in its later deformational stage. The stable temperature distribution profile in the ASB stems from little plastic work done by the flow shear stress acting on the ASB. The present analysis is carried out in the adiabatic condition and the assumption that plastic work rate and heat conduction reach equilibrium is not necessary in obtaining a stable ASB width. In the adiabatic condition,

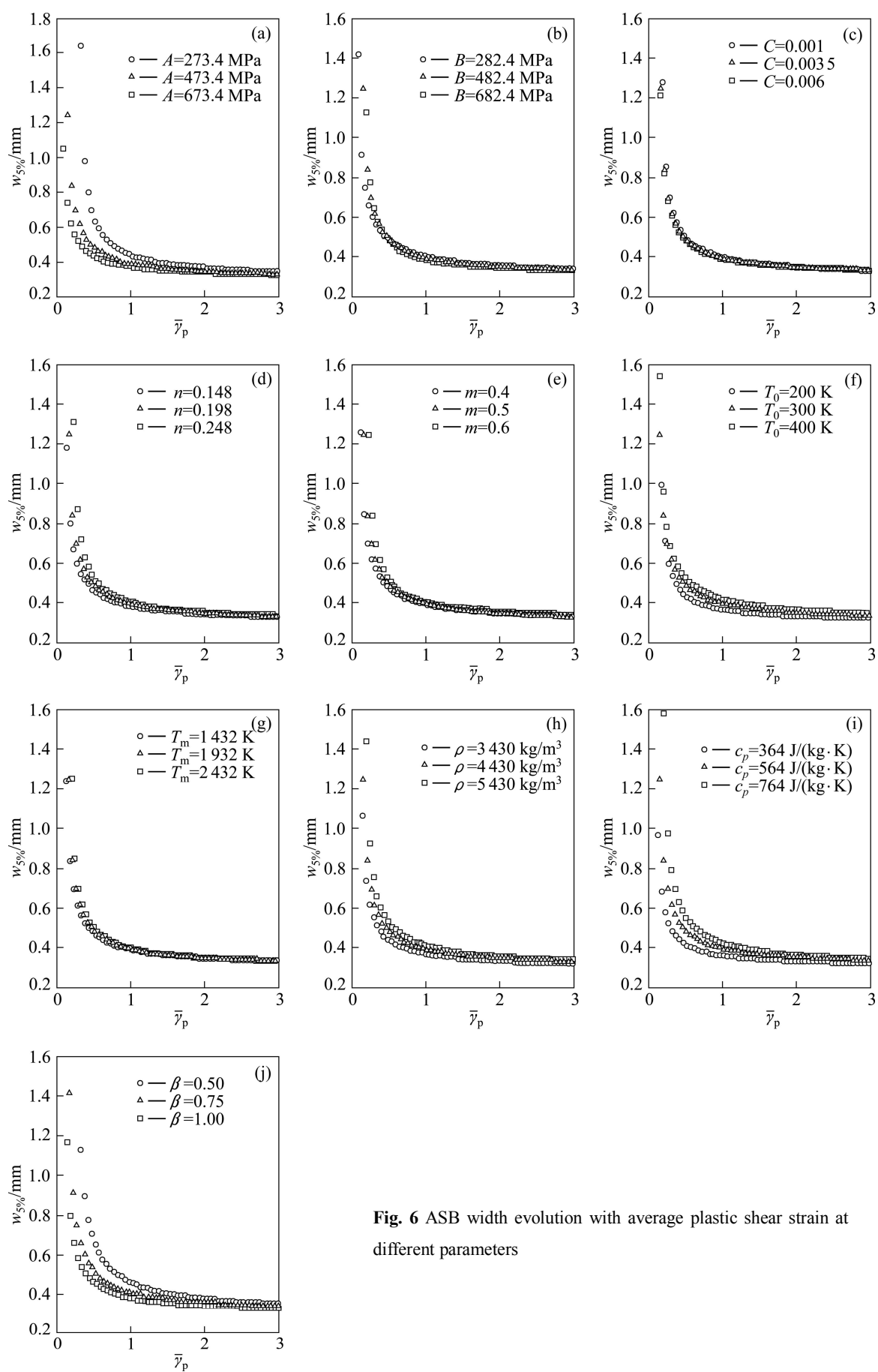


Fig. 6 ASB width evolution with average plastic shear strain at different parameters

the model proposed by DODD and BAI [5] cannot be applied. In addition, the assumed equilibrium condition is satisfied difficultly at high strain rates. The model proposed by DODD and BAI [5] only can predict the stable ASB width, while the present model proposed can predict the overall evolution process that an ASB becomes thin until a stable value is reached.

4 Conclusions

1) To assess the critical plastic shear strain corresponding to the onset of the ASB for Ti-6Al-4V, the experimental data for a flow line in the ASB obtained by LIAO and DUFFY are fitted by use of the curve-fitting least squares method and the theoretical expression for the local plastic shear strain distribution proposed by WANG. It is found that the proposed expression well describes the non-linear deformation characteristics of the flow line in the ASB.

2) Some parameters in the JOHNSON-COOK model are back-calculated at different critical plastic shear strains. The influence of 10 parameters on the adiabatic shear sensitivity is studied. The present model can predict the ASB width evolution process. The predicted ASB width decreases with straining until a stable value is reached. The famous model proposed by DODD and BAI only can predict a final stable value.

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基于梯度依赖的 JOHNSON-COOK 模型的 韧性绝热剪切敏感性

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摘 要: 以基于梯度塑性理论提出的绝热剪切带内部的局部塑性剪切变形分布的理论表达式为基础, 研究 10 个参数对绝热剪切敏感性的影响。对 LIAO 及 DUFFY 给出的 Ti-6Al-4V 绝热剪切带内部的 1 条流线的实验结果进行最小二乘曲线拟合。估算绝热剪切带宽度取值不同时的临界塑性剪切应变。理论曲线很好地反映了绝热剪切带内部流线的非线性变形特征。利用不同的临界塑性剪切应变值反算了 JOHNSON-COOK 模型中的一些参数。研究发现, 绝热剪切敏感性随着初始静态屈服应力、功热转化因子和应变率参数的降低而降低, 这与密度、热容、环境温度及应变硬化指数的影响刚好相反。所提出的模型可以预测绝热剪切带的宽度由高至低的演变过程, 直到达到一个稳定值, 这一点 DODD 和 BAI 模型做不到。

关键词: 绝热剪切带; Ti-6Al-4V; JOHNSON-COOK 模型; 宽度; 梯度塑性理论; 局部变形分布

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