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Sensitivity analysis on globularized fraction of α lamellae in titanium alloys

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Abstract: The sensitivity analysis functions on globularized fraction of α lamellae were established using a physically-based microstructure model and gradient method. These functions were applied to the sensitivity analysis on globularized fraction of α lamellae in TC17 alloy. The material constants in these functions are determined using the genetic algorithm-based objective optimization technique. The globularized fraction of α lamellae during isothermal compression of TC17 alloy was quantitatively analyzed based on scanning electron microscopy (SEM) observation. The results show that α lamellae mostly change to equiaxed α grains at a deformation temperature of 1083 K, a strain rate of 0.01 s^{-1} and a strain of 1.2. The globularized fraction decreases with increasing strain rate because lower strain rate provides enough time for the spheroidization. The effect of deformation temperature on the globularized fraction is controlled by the strain rate. And, the predicted derivations of globularized fraction with respect to processing parameters show good agreement with the experimental values.

Key words: titanium alloys; sensitivity analysis; globularized fraction; genetic algorithm

1 Introduction

It is well known that the mechanical properties of titanium alloys are strongly affected by the microstructural characteristics (e.g., the volume fraction of various phases, grain size, and grain misorientation). During high temperature deformation, the alloys undergo significant microstructural changes including dynamic recrystallization or recovery and grain growth. The microstructural evolution is very sensitive to the processing parameters, and strongly affects the mechanical properties [1–4]. Therefore, it is very important to obtain the law of the microstructural evolution, and make the sensitivity analysis of microstructure with respect to processing parameters so as to control the microstructure.

In recent years, extensive efforts were paid to simulating the microstructural evolution and making the sensitivity analysis of microstructure for a wide range of metals and alloys [5]. KOHAR et al [6] established a new multi-scale framework which incorporated the microstructural evolution, and stimulated texture

evolution in single crystal and polycrystal FCC material. FULLWOOD et al [7] got the microstructure–property relations in Fourier form, leading to an efficient spectral framework of sensitive analysis and designing for microstructure. McDOWELL and DUNNE [8] developed the microstructure-sensitive model and facilitated parametric design exploration in searching for morphologies and/or compositions that modified fatigue resistance. Much prior work showed that the sensitivity analysis of microstructure was made based on the empirical microstructure model and finite element simulation. Although these sensitivity analyses could provide a optimization scheme to control the microstructure and design the processing parameters during deformation, the empirical microstructure model cannot reveal the plastic deformation mechanisms of metals and alloys. It is well known that the physically-based microstructure model is not only highly reliable in predicting microstructure, but it can also reveal the plastic deformation mechanisms.

Therefore, the object of the present work is to make the sensitivity analysis of microstructure with respect to processing parameters with the aid of physically-based

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microstructure model of titanium alloys. For this purpose, the gradient functions are expressed as the sensitivities of strain, deformation temperature and strain rate with respect to globularized fraction of α lamellae. Applying the sensitivity analysis to high temperature deformation process of TC17 alloy, the sensitivity domains of the globularized fraction are confirmed to control the microstructure in the present work.

2 Experimental

The 50 mm-in-diameter billets of TC17 alloy had chemical composition (wt.%) of 5.12 Al, 2.03 Sn, 2.10 Zr, 4.04 Mo, 3.94 Cr, 0.10 Fe, 0.12 C, 0.007 N, 0.007 H, 0.12 O, and balanced Ti. TC17 alloy containing lamellar α microstructure was compressed at the deformation temperatures of 1043, 1063, 1083, 1103 and 1123 K, the strain rates of 0.01, 0.1, 1.0 and 5.0 s^{-1} , and the strains of 0.2, 0.5, 0.7, 0.9 and 1.2. A thermocouple was welded on the surface of the specimens to measure the temperature during deformation, and graphite powder was put between the specimens and the anvils to reduce die friction. After isothermal compression, the specimens were water-quenched to room temperature. To observe deformed microstructure, the specimens were mechanically polished and chemically etched in a solution of 10 mL HF, 15 mL HNO_3 and 75 mL H_2O . A Zeiss SUPRA 55 SEM operating at 20 kV was used to examine the microstructural evolution of lamellar α microstructure in TC17 alloy. The SEM micrographs were taken along the direction perpendicular to the compression axis. The globularized fraction of lamellar α microstructure in TC17 alloy was measured using quantitative metallography image analysis software (Image-Pro Plus 6.0), and it was calculated by the average value of twenty visual fields.

3 Sensitivity analysis of globularized fraction

In the present work, it is assumed that the globularized fraction follows an Avrami type equation [9,10]:

$$X(t) = 1 - \exp\left(-\frac{4\pi}{3} N_V (Gt)^3\right) \quad (1)$$

where $X(t)$ is the globularized fraction at time t ; G is the growth rate of globular grain (m/s); N_V is the number of globular grain per unit volume, and it is given by [10,11]

$$N_V = \frac{4S_0(1 - \exp(-\phi t))}{\pi d^2} \quad (2)$$

where S_0 is the initial surface area per unit volume (μm^{-1}), t is the time (s), ϕ is the material constant (s^{-1}),

and d is the size of globular grain (μm).

G is expressed as [12,13]

$$G = \frac{M_{\text{gb}}\sigma_{\text{surf}}}{d} = \frac{b\delta D_0\sigma_{\text{surf}} \exp(-\frac{Q_{\text{pd}}}{RT})}{kTd} \quad (3)$$

where M_{gb} is the grain boundary mobility ($\text{m}^4\text{J}^{-1}\text{s}^{-1}$), σ_{surf} is the grain boundary energy per unit area (J/m^2), b is the magnitude of Burgers vector ($2.86 \times 10^{-10} \text{ m}$), δ is the characteristic grain boundary thickness (m), D_0 is the self-diffusion constant, k is Boltzmann's constant ($1.381 \times 10^{-23} \text{ J/K}$), and Q_{pd} is the boundary diffusion activation energy (kJ/mol).

If the deformation is performed at a constant strain rate, there is

$$t = \varepsilon/\dot{\varepsilon} \quad (4)$$

where ε is the strain, and $\dot{\varepsilon}$ is the strain rate (s^{-1}).

By substituting Eqs. (2)–(4) into Eq. (1), the globularized fraction X is written as

$$X = 1 - \exp(-\gamma_1 \varepsilon^3 (1 - \exp(-\gamma_2 \varepsilon))) \quad (5)$$

$$\gamma_1 = \frac{16S_0(b\delta D_0\sigma_{\text{surf}})^3}{3(kT\dot{\varepsilon})^3 d^5} \exp(-\frac{3Q_{\text{pd}}}{RT}), \quad \gamma_2 = \phi\dot{\varepsilon}^{-1}$$

According to the transformation kinetics law, the globularized fraction X is modified by

$$X = 1 - \exp(-\gamma_1 \varepsilon^n (1 - \exp(-\gamma_2 \varepsilon))) \quad (6)$$

where n is the kinetics exponent; γ_1 and γ_2 are the material constants.

In addition, the Zener–Hollomon parameter is introduced into Eq. (6) to describe the effect of deformation temperature and strain rate on the globularized fraction X . Finally, the globularized fraction X is written in the following form:

$$X = 1 - \exp(-f(z)\varepsilon^n (1 - \exp(-\gamma_3 \varepsilon))) \quad (7)$$

where $f(z) = pZ^m$; $Z = \dot{\varepsilon} \exp[Q/(RT)]$; p , m and γ_3 are the material constants; Q is the apparent activation energy for deformation (kJ/mol), R is the gas constant ($8.3145 \text{ J}\cdot\text{mol}^{-1}\cdot\text{K}^{-1}$); T is the absolute deformation temperature (K).

By differentiating Eq. (7), the gradients of the globularized fraction X with respect to the strain, the deformation temperature and the $\lg(\text{strain rate})$ are calculated as

$$\frac{\partial X}{\partial \varepsilon} = \exp(-p(\dot{\varepsilon} \exp(\frac{Q}{RT}))^m \varepsilon^n (1 - \exp(-\gamma_3 \varepsilon))) np \cdot$$

$$(\dot{\varepsilon} \exp(\frac{Q}{RT}))^m \varepsilon^{n-1} (1 + (\frac{\varepsilon}{n} - 1) \exp(-\gamma_3 \varepsilon)) \quad (8)$$

$$\frac{\partial X}{\partial T} = -p\varepsilon^n(1-\exp(-\gamma_3\varepsilon))\dot{\varepsilon}^m \frac{mQ}{R} \exp(-p\varepsilon^n \cdot (1-\exp(-\gamma_3\varepsilon))\dot{\varepsilon}^m \exp(\frac{mQ}{RT}))T^{-2} \exp(\frac{mQ}{RT}) \quad (9)$$

$$\frac{\partial X}{\partial(\lg\dot{\varepsilon})} = \ln 10 * m p \varepsilon^n (1-\exp(-\gamma_3\varepsilon)) \exp(\frac{mQ}{RT}) 10^{m\lg\dot{\varepsilon}} \cdot \exp(-p\varepsilon^n(1-\exp(-\gamma_3\varepsilon))\exp(\frac{mQ}{RT})10^{m\lg\dot{\varepsilon}}) \quad (10)$$

From Eqs. (8)–(10), it is seen that these gradient functions can be used for the sensitivity analysis of globularized fraction during high temperature deformation of titanium alloys.

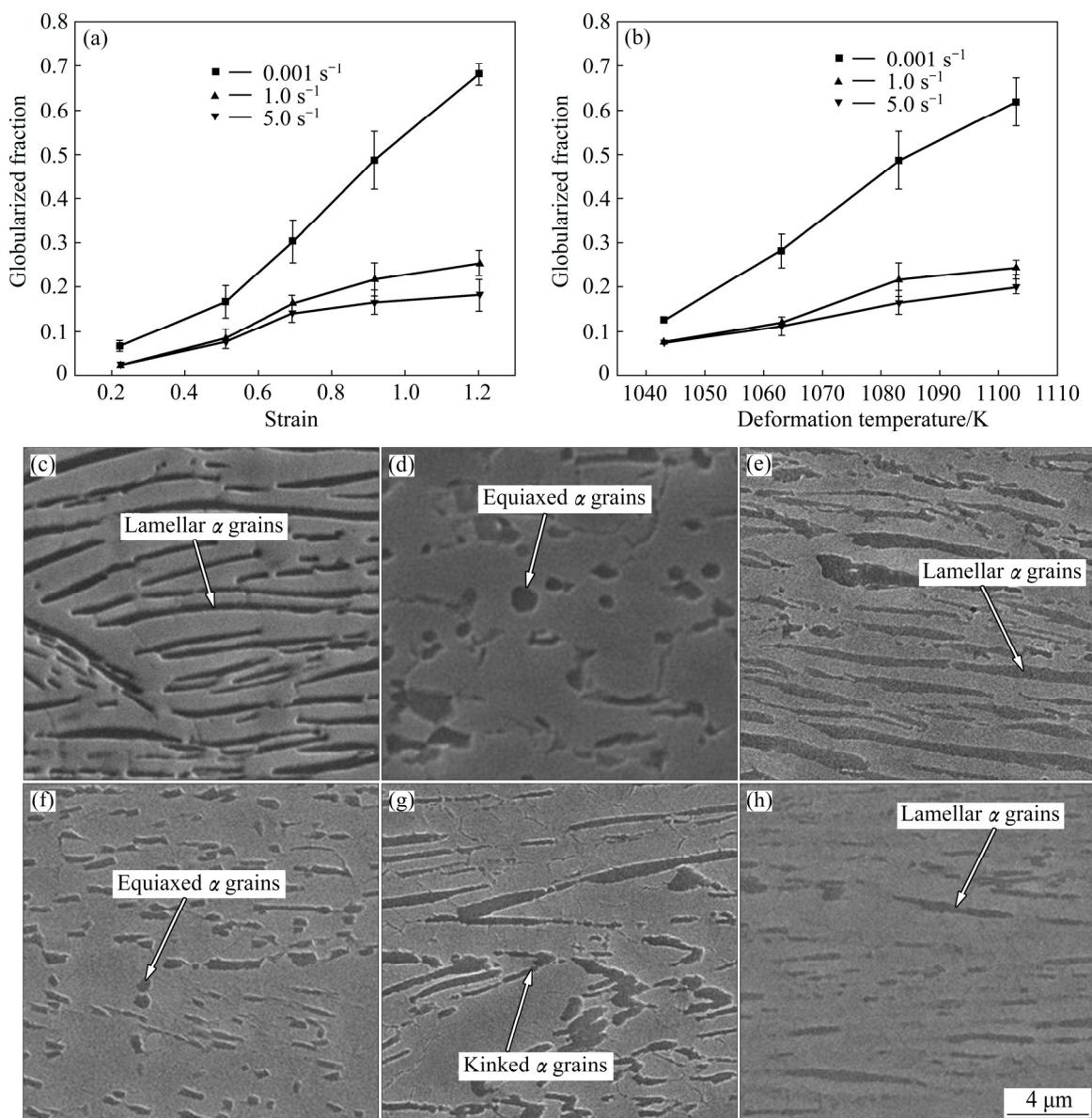


Fig. 1 Effect of processing parameters on globularized fraction of lamellar α microstructure in TC17 alloy (a, b) and SEM micrographs with different processing parameters (c–h): (a) $T=1083$ K; (b) $\varepsilon=0.9$; (c) $T=1083$ K, $\dot{\varepsilon}=0.01$ s⁻¹, $\varepsilon=0.2$; (d) $T=1083$ K, $\dot{\varepsilon}=0.01$ s⁻¹, $\varepsilon=1.2$; (e) $T=1043$ K, $\dot{\varepsilon}=0.01$ s⁻¹, $\varepsilon=0.9$; (f) $T=1103$ K, $\dot{\varepsilon}=0.01$ s⁻¹, $\varepsilon=0.9$; (g) $T=1103$ K, $\dot{\varepsilon}=1.0$ s⁻¹, $\varepsilon=0.9$; (h) $T=1103$ K, $\dot{\varepsilon}=5.0$ s⁻¹, $\varepsilon=0.9$

4 Applying functions of sensitivity analysis to lamellar α microstructure in TC17 alloy

4.1 Effect of processing parameters on globularized fraction

The globularization behavior of lamellar α microstructure in TC17 alloy is quantified with the aid of SEM observations and quantitative metallography image analysis software. Figure 1 shows the globularized fraction as a function of strain, deformation temperature and strain rate. At a deformation temperature of 1083 K and a strain rate of 0.01 s⁻¹, the globularized fraction significantly increases from 0.07 to 0.68 when the strain

increases from 0.2 to 1.2. The SEM micrographs at the strains of 0.2 and 1.2 are shown in Fig. 1(c) and Fig. 1(d), respectively. In Fig. 1(c), the parallel α lamellae are present in this alloy, without the formation of significant equiaxed α grains. At a strain of 1.2, almost fully equiaxed microstructure that contains only a small amount of unglobularized α lamellae is provided in Fig. 1(d), suggesting the fragmentation and globularization of α lamellae via groove/boundary splitting under this condition [14–16]. At the strain rates of 1.0 and 5.0 s^{-1} , the globularized fraction of TC17 alloy also shows a similar increase with increasing strain, but the increasing extent at the strain rates of 1.0 and 5.0 s^{-1} decreases, as illustrated in Fig. 1(a).

Figure 1(b) shows the effect of deformation temperature on the globularized fraction of lamellar α microstructure in TC17 alloy. At a strain rate of 0.01 s^{-1} and a strain of 0.9, the globularized fraction significantly increases from 0.12 to 0.62 when the deformation temperature increases from 1043 to 1103 K. Figures 1(e) and (f) show the SEM micrographs at the deformation temperatures of 1043 and 1103 K, respectively. At 1043 K (Fig. 1(e)), the microstructure is comprised of the parallel α lamellae and β phase. When the deformation temperature is up to 1103 K (Fig. 1(f)), the microstructure is comprised of equiaxed α grains, a small amount of unglobularized α lamellae and β phase, consistent with the quantitative results above. This phenomenon demonstrates that the deformation temperature has a significant effect on the globularized fraction of TC17 alloy at a strain rate of 0.01 s^{-1} and a strain of 0.9. However, the quantitative results of the globularized fraction reveal a weak dependence on the deformation temperature at the strain rate of 1.0 and 5.0 s^{-1} . Main reason is that lower strain rate provides enough time for the spheroidization of lamellar α microstructure in TC17 alloy, and higher deformation temperature is beneficial for the diffusion of solutes. So, enough time and high solute diffusion promote a significant increase in globularized fraction of lamellar α microstructure in TC17 alloy.

In addition, it is seen that the globularized fraction of lamellar α microstructure in TC17 alloy decreases with increasing strain rate, as illustrated in Figs. 1(a) and (b). This phenomenon can also be confirmed by the SEM micrographs in Figs. 1(f)–(h). It is seen in Figs. 1(f)–(h) that equiaxed α grains decrease with increasing strain rate. Figure 1(h) shows that a few bent and kinked α lamellae are present under this condition, and α lamellae are parallel along the direction perpendicular to the compression axis. At a strain rate of 5.0 s^{-1} (Fig. 1(h)), only parallel α lamellae are observed. This is possibly due to lower strain rate providing enough time for the spheroidization. The above results and analyses indicate

that the globularization behavior of lamellar α microstructure in TC17 alloy is highly dependent on the strain and strain rate during deformation. However, the effect of deformation temperature on the globularization behavior of lamellar α microstructure in TC17 alloy is controlled by the strain rate.

4.2 Identification of material constants in gradient functions

The measurements of the globularized fraction as a function of strain, deformation temperature and strain rate of lamellar α microstructure in TC17 alloy are shown in Fig. 1. Based on these measured values, the derivatives of the fraction globularized with respect to processing parameters (strain, deformation temperature, $\lg(\text{strain rate})$) can be obtained as the experimental data. Using conventional optimization method, it is difficult to determine the material constants in the gradient functions (Eqs. (8)–(10)). The material constants in these gradient functions are determined using the GA-based objective optimization technique [17,18]. The objective functions are expressed as

$$f_1(x) = \sum_{k=1}^{l_1} \sum_{j=1}^{n_1} \sum_{i=1}^{m_1} w_{ijk} \left[\left(\left(\frac{\partial X}{\partial \varepsilon} \right)_i^c \right)_j \right]_k - \left(\left(\left(\frac{\partial X}{\partial \varepsilon} \right)_i^e \right)_j \right)_k \right]^2 \quad (11)$$

$$f_2(x) = \sum_{p=1}^{l_2} \sum_{q=1}^{n_2} \sum_{h=1}^{m_2} w_{pqh} \left[\left(\left(\frac{\partial X}{\partial T} \right)_p^c \right)_q \right]_h - \left(\left(\left(\frac{\partial X}{\partial T} \right)_p^e \right)_q \right)_h \right]^2 \quad (12)$$

$$f_3(x) = \sum_{y=1}^{l_3} \sum_{s=1}^{n_3} \sum_{z=1}^{m_3} w_{ysz} \cdot \left[\left(\left(\frac{\partial X}{\partial \lg \dot{\varepsilon}} \right)_y^c \right)_s \right]_z - \left(\left(\left(\frac{\partial X}{\partial \lg \dot{\varepsilon}} \right)_y^e \right)_s \right)_z \right]^2 \quad (13)$$

where ε is the strain; T is the deformation temperature (K); $\dot{\varepsilon}$ is the strain rate (s^{-1}); $f_1(x)$, $f_2(x)$ and $f_3(x)$ are the residuals for the derivatives of the globularized fraction with respect to strain, deformation temperature and $\lg(\text{strain rate})$, respectively; X is the globularized fraction; $\left(\left(\frac{\partial X}{\partial \varepsilon} \right)_i^c \right)_j$ and $\left(\left(\frac{\partial X}{\partial \varepsilon} \right)_i^e \right)_j$ are the predicted and experimental derivatives of the globularized fraction with respect to strain at strain i , strain rate j and deformation temperature k ; m_1 is the number of strain; n_1 is the number of strain rate; l_1 is the number of deformation temperature; w_{ijk} is the weight coefficient. Similarly, $\left(\left(\frac{\partial X}{\partial T} \right)_p^c \right)_q$ and $\left(\left(\frac{\partial X}{\partial T} \right)_p^e \right)_q$ are the predicted and experimental derivatives of the globularized fraction with respect to deformation

temperature at time p , strain rate q and deformation temperature h ; m_2 is the number of strain; n_2 is the number of strain rate; l_2 is the number of deformation temperature; w_{pgh} is the weight coefficient. $((\partial X / \partial \lg \dot{\varepsilon})_y^c)_z$ and $((\partial X / \partial \lg \dot{\varepsilon})_y^e)_z$ are the predicted and experimental derivatives of the globularized fraction with respect to $\lg(\text{strain rate})$ at time y , strain rate s and deformation temperature z ; m_3 is the number of strain; n_3 is the number of strain rate; l_3 is the number of deformation temperature; w_{ysz} is the weight coefficient. Finally, the material constants in Eqs. (8)–(10) are listed in Table 1.

Table 1 Material constants in gradient functions of lamellar α microstructure in TC17 alloy

p	$Q/(\text{kJ} \cdot \text{mol}^{-1})$	m	n	γ_3
0.85	555.03	0.17	1.27	13.08

4.3 Comparison of predicted data with experimental data

Using the material constants in Table 1, the derivatives of the globularized fraction with respect to

processing parameters (strain, deformation temperature, and $\lg(\text{strain rate})$) are predicted. Figure 2 shows the comparison of the predicted with the experimental $\partial X / \partial \varepsilon$. It is seen that the average relative differences for the sample data and the non-sample data are 8.8% and 11.8%, respectively. Figure 3 shows the comparison of the predicted with the experimental $\partial X / \partial T$. It is seen that the average relative differences for the sample data and the non-sample data are 9.2% and 10.4%, respectively. Figure 4 shows the comparison of the predicted with the experimental $\partial X / \partial \lg \dot{\varepsilon}$. It is seen that the average relative differences for sample data and non-sample data are 8.1% and 11.3%, respectively. The above results indicate that these gradient functions with a high prediction precision can be used to predict the derivatives of the globularized fraction with respect to processing parameters in TC17 alloy containing lamellar α microstructure.

4.4 Sensitivity domains of globularized fraction in TC17 alloy

Based on the predicted derivatives of the globularized fraction with respect to processing

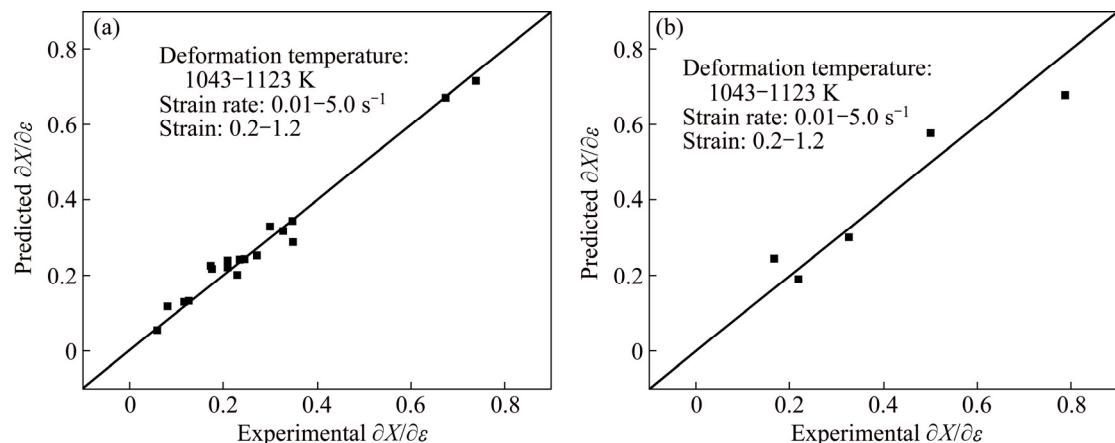


Fig. 2 Comparison of predicted with experimental $\partial X / \partial \varepsilon$: (a) Sample data; (b) Non-sample data

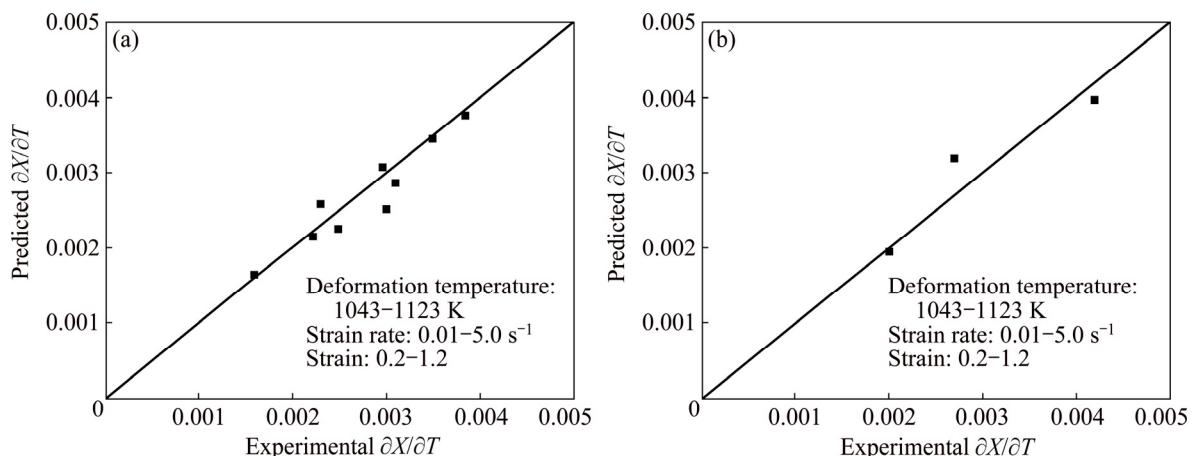


Fig. 3 Comparison of predicted with experimental $\partial X / \partial T$: (a) Sample data; (b) Non-sample data

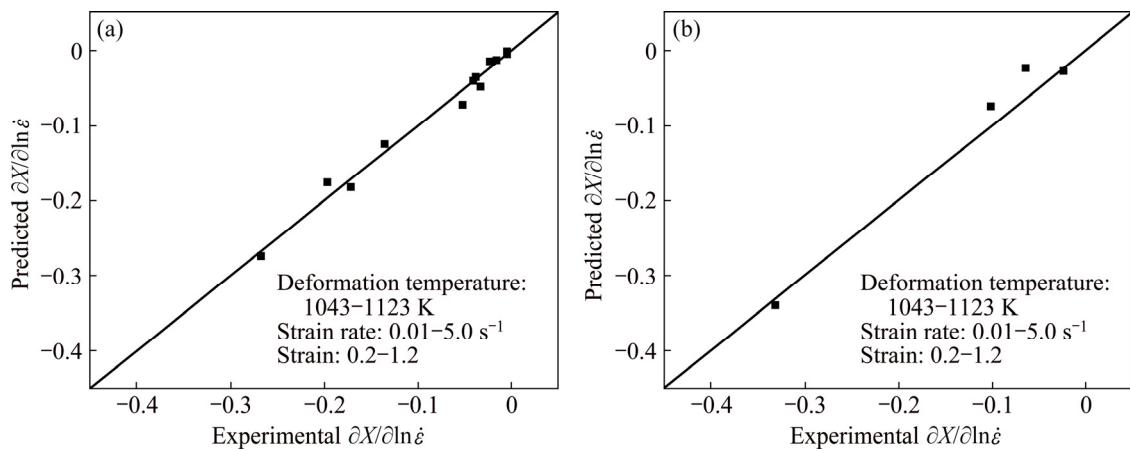


Fig. 4 Comparison of predicted with experimental $\partial X / \partial \ln \dot{\varepsilon}$: (a) Sample data; (b) Non-sample data

parameters, the sensitivity domains of lamellar α microstructure in TC17 alloy are confirmed. Figure 5 shows the predicted curves for the derivatives of globularized fraction with respect to processing parameters, implying that the predicted values show good agreement with the experimental results. In Fig. 5(a), it is seen that the derivative of globularized fraction with respect to strain decreases with increasing strain after a peak value. The possible reason is that large strain produces very high dislocation accumulation and local dislocation tangles in early stage of deformation, which promotes the formation of low-angle boundaries (LABs) in α lamella interiors and finally results in the occurrence of the fragmentation and globularization of α lamellae via groove/boundary splitting. Thus, the derivative of globularized fraction with respect to strain firstly increases with increasing strain. However, continuous fragmentation and globularization of α lamellae become more difficult after a certain globularized degree due to the fact that most plastic deformation energy is used for the grain boundary sliding of equiaxed α grains in the late stage of deformation [19]. So, the derivative of globularized fraction with respect to strain decreases with increasing strain after a peak value. In addition, it is seen that the derivative value of globularized fraction with respect to strain at a deformation temperature of 1083 K and a strain rate of 0.01 s^{-1} is the largest when the strain is in the range of 0.7–1.2, implying that the globularized fraction of lamellar α microstructure in TC17 alloy is the most sensitive to the strain at this deformation temperature and strain rate. The above analysis indicates that the globularization behavior of lamellar α microstructure in TC17 alloy is very noticeable if the deformation is performed in these sensitivity domains.

In Fig. 5(b), it is seen that the derivative of globularized fraction with respect to deformation

temperature has a general increasing tendency. This is possibly attributed to the fact that the globularization process is greatly affected by the diffusion-controlled interface migration. The interface migration that is a thermally-activated process is influenced by the deformation temperature. Thus, the derivative of globularized fraction with respect to deformation temperature has a general increasing tendency with increasing deformation temperature. In addition, the derivative value of globularized fraction with respect to deformation temperature at a strain of 0.9 and a strain rate of 0.1 s^{-1} is the largest when the deformation temperature is in the range of 1063–1103 K, implying that the globularized fraction of lamellar α microstructure in TC17 alloy is the most sensitive to the deformation temperature at this strain and strain rate. The above analysis indicates that the globularization behavior of lamellar α microstructure in TC17 alloy is very noticeable if the deformation is performed in these sensitivity domains.

In Fig. 5(c), it is seen that the absolute derivative of globularized fraction with respect to $\lg(\text{strain rate})$ decreases with increasing strain rate. Main reason is that lower strain rate provides enough time for the spheroidization of lamellar α microstructure in TC17 alloy. In addition, the absolute derivative value of globularized fraction with respect to $\lg(\text{strain rate})$ at a strain of 0.9 and a deformation temperature of 1083–1103 K is the largest when the strain rate is in the range of $0.01\text{--}0.1 \text{ s}^{-1}$, implying that the globularized fraction of lamellar α microstructure in TC17 alloy is the most sensitive to the strain rate at this strain and in this deformation temperature range. The above analysis indicates that the globularization behavior of TC17 alloy containing lamellar α microstructure is very noticeable if the deformation is performed in these sensitivity domains.

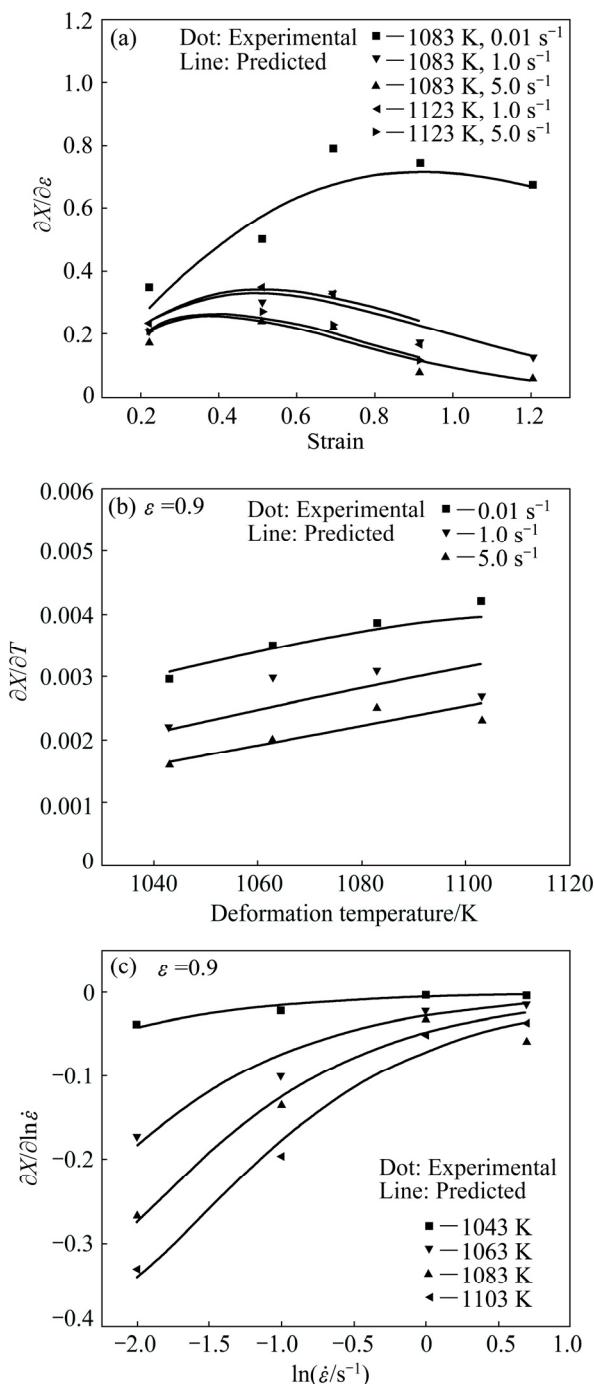


Fig. 5 Predicted curves for derivatives of globularized fraction with respect to processing parameters: (a) Strain; (b) Deformation temperature; (c) lg(strain rate)

5 Conclusions

(1) When the strain increases to 1.2, the α lamellae mostly change to equiaxed α grains at 1083 K and 0.01 s⁻¹, suggesting the fragmentation and globularization of α lamellae via groove/boundary splitting.

(2) The globularized fraction decreases with increasing strain rate because lower strain rate provides

enough time for the spheroidization. The effect of deformation temperature on the fraction globularized is controlled by strain rate.

(3) The gradient functions of the globularized fraction with respect to the strain, the deformation temperature and the $\lg(\text{strain rate})$ are obtained. The results indicate that these gradient functions with a high prediction precision can be used to predict the derivatives of the globularized fraction with respect to processing parameters in TC17 alloy containing lamellar α microstructure.

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钛合金中 α 片层球化分数敏感性分析

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摘要: 基于机理型微观组织模型与梯度算法建立钛合金中 α 片层球化分数敏感性分析函数, 将该函数应用于 TC17 合金中片层组织球化分数的敏感性分析。基于扫描电镜观测结果定量分析 TC17 合金等温压缩过程中工艺参数对球化分数的影响规律, 并采用遗传算法优化其敏感性分析函数的材料参数。结果表明: 当变形温度为 1083 K、应变速率为 0.01 s^{-1} 、应变为 1.2 时, TC17 合金片层 α 组织几乎完全转变为等轴 α 晶粒; 随着应变速率的增加, 片层 α 球化分数减少, 这主要因为较低的应变速率为动态球化提供了足够的时间; 而变形温度对片层 α 球化分数的影响受应变速率控制。此外, TC17 合金片层 α 球化分数关于应变、变形温度和对数应变速率微分的预测结果与试验结果相吻合。

关键词: 钛合金; 敏感性分析; 球化分数; 遗传算法

(Edited by Bing YANG)