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Electronic structures and physical properties of pure Cr, $\,$ Mo and $\,W^{^{\odot}}$

XIE Yourqing(谢佑卿), DENG Yong-ping(邓永平), LIU Xin-bi(刘心笔) (School of Materials Science and Engineering, Central South University, Changsha 410083, China)

Abstract: Using the one atom theory, the electronic structures of pure Cr, Mo and W with bcc structure were determined respectively as: $[Ar](3d_c)^{3.32}(3d_n)^{2.26}(4s_c)^{0.25}(4s_f)^{0.17}$, $[Kr](4d_c)^{4.23}(4d_n)^{1.48}(5s_c)^{0.02}(5s_f)^{0.27}$ and $[Xe](5d_c)^{5.16}(6s_c)^{0.25}(6s_f)^{0.59}$. The electronic structures of these metals with hcp and fcc structures and liquid state were also studied. According to their electronic structures, the relationship between the electronic structure and crystalline structure was explained qualitatively and the relationship between the difference of mechanical properties and transport properties of pure Cr, Mo and W with bcc structure and their electronic structures was also explained qualitatively; the lattice comstants, binding energy, potential curves, elasticities and the temperature dependence of the linear thermal expansion coefficient of bcc Cr, bcc Mo and bcc W were calculated quantitatively.

Key words: Cr; Mo; W; electronic structure; crystalline structure; lattice constant; binding energy; mechanical property; transport property

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1 INTRODUCTION

In the framework of the metallic materials systematic science^[1], there are three scientific theory systems: pure metal systematic science, alloys physics and chemistry, and alloys statistical thermodynamics. The core of the pure metal systemætic science is the one atom theory (OA)^[2-5], and that of the alloys physics and chemistry and alloys statistical thermodynamics is the characteristic crystal theory (CC). In order to make the materials design step from empirical to scientific under the guidance of the metallic materials systematic science, the first step is to establish a pure metal data base according to OA theory. Using the auto-search-solution OA theory program developed by the authors, the calculation for electronic structures and physical properties of pure metal Cr, Mo, W, Pt, Ru, etc have been completed, and that of Cr, Mo, W are introduced in this paper.

2 BASIC ATOM STATES OF Cr, Mo AND W

In the OA theory, the electronic structures of pure metals are described by the quasi-electron-occupation (QEO) number of the one atom state Ψ_a which is made up of some basic atom states Φ_k :

$$\psi_{\mathbf{a}} = \sum_{k} c_k \Phi_k \tag{1}$$

At the outer shell of Cr, Mo, W atom, there are covalent electrons n_c , near-free electrons n_f and no-

bond electrons n_n . In each basic atom state, the distribution of electrons follows the Pauli exclusion principle.

The characteristic properties (lattice constant a and binding energy E_c) of the pseudo-crystals formed by atoms in each kind of basic atom state of Cr, Mo, W metal can be obtained using a series of established expressions^[3]. The results of Cr metal are listed in Table 1.

3 ELECTRONIC STRUCTURES OF Cr, Mo, W

3. 1 Electronic structures of bcc Cr, bcc Mo and bcc W

In our research, firstly we take the two characteristic properties (a and E_c) as the criterion and perform the three-state hybridization to all the 17 basic atom states with the step length of 0. 01. Secondly, after a systematic analysis mainly taking the lattice constants and binding energies into consideration, the three-state combination of bcc Cr, bcc Mo and bcc W are determined as follows:

$$c_2 = 0.04$$
, $c_{12} = 0.79$, $c_{14} = 0.17$; $c_4 = 0.27$, $c_{12} = 0.72$, $c_{13} = 0.01$; $c_1 = 0.17$, $c_3 = 0.25$, $c_6 = 0.58$.

Then the electronic structures of bcc Cr, bcc Mo and bcc W are determined as

$$\begin{array}{ll} \Psi_{\!a}\!(\;bcc\;Cr) = & [\;Ar]\;(\;3d_c\,)^{\;3.\;32}\,(\;3d_n\,)^{\;2.\;26}\,(\;4s_c\,)^{\;0.\;25}\,\bullet \\ (\;4s_f)^{\;0.\;17}; \end{array}$$

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$$\begin{array}{lll} \Psi_{a}\;(\;bcc\,M\,o\,)\;=\;\;[\;Kr]\;\;(\;4d_{c}\,)^{\;4.\;23}\;(\;4d_{n}\,)^{\;1.\;48}\;\bullet\\ (\;5s_{c})^{\;0.\;02}(\;5s_{f})^{\;0.\;27};\\ \Psi_{a}(\;bcc\,W\,)\;=\;\;[\;Xe]\;(\;5d_{c})^{\;5.\;16}(\;6s_{c})^{\;0.\;25}(\;6s_{f})^{\;0.\;59}. \end{array}$$

This process of determining electronic structure is multiple properties determining state method in OA theory^[6]. The atom state parameters, bond parameters and characteristic properties of bcc Cr, bcc Mo and bcc W crystal are listed in Table 2, where r_1 , r_2 , r_3 and n_1 , n_2 , n_3 denote covalent bond length and covalent electron pair number, respectively.

3. 2 Electronic structures of hcp Cr, fcc Cr, hcp Mo, fcc Mo, hcp W and fcc W

3. 2. 1 Binding energies of hcp Cr, fcc Cr, hcp Mo, fcc Mo, hcp W and fcc W

Under the isopiestic condition, the Gibbs energy of pure metal is a function of specific heat capacity $C_p(T)$ and volume V(T):

$$G(T) = E_{c} + \int_{0}^{T} C_{p}(T) dT + p \int_{0}^{T} V(T) dT - T \int_{0}^{T} C_{p}(T) / T dT$$
 (2)

where p is pressure and T is temperature.

The relative Gibbs energy is presented by a polynomial of T in SGTE(Scientific Group Thermodata, Europe) database^[7]:

 $G'(T) = a + bT + cT \ln T + \sum k_n T^n$ (3) where G'(T) is the relative Gibbs energy, and it is not equal to G(T). At 0 K, there exists

 $G(T) - G'(T) = E_c - a \tag{4}$

Table 1 Basic atom states and corresponding pseudo-crystal characteristic properties of Cr

7	Electronic structure	area arra corr	a/nm	seudo erystar	$E_{c}/\left(kJ^{\bullet}mol^{-1}\right)$		
k	in outer shell	bec	hep	fee	bcc	hcp	fee
1	$(3d_c)^{4}(3d_n)^{0}(4s_c)^{0}(4s_f)^{2}$	0. 310 08	0. 275 53	0.389 68	371.75	367. 31	362.04
2	$(3d_c)^{4}(3d_n)^{0}(4s_c)^{2}(4s_f)^{0}$	0. 298 15	0. 265 00	0.374 79	650. 78	645. 17	644.96
3	$(3d_c)^{4}(3d_n)^{0}(4s_c)^{1}(4s_f)^{1}$	0. 303 51	0. 269 74	0. 381 49	515. 18	510. 53	508.23
4	$(3d_c)^{5}(3d_n)^{0}(4s_c)^{0}(4s_f)^{1}$	0. 284 71	0. 253 13	0.358 01	502. 66	498. 52	496.01
5	$(3d_c)^{5}(3d_n)^{0}(4s_c)^{1}(4s_f)^{0}$	0. 279 36	0. 248 40	0.35133	706. 47	701. 07	700.69
6	$(3d_c)^{6}(3d_n)^{0}(4s_c)^{0}(4s_f)^{0}$	0. 260 62	0. 231 82	0.327 90	689. 69	685. 05	684.46
7	$(3d_c)^{2}(3d_n)^{2}(4s_c)^{0}(4s_f)^{2}$	0.358 80	0. 318 88	0.450 98	176. 27	174. 39	168.03
8	$(3d_c)^{2}(3d_n)^{2}(4s_c)^{2}(4s_f)^{0}$	0. 328 94	0. 292 16	0.413 19	375. 51	371. 83	371.75
9	$(3d_c)^{2}(3d_n)^{2}(4s_c)^{1}(4s_f)^{1}$	0. 337 43	0. 306 42	0.433 37	271. 54	262. 46	260.04
10	$(3d_c)^{3}(3d_n)^{2}(4s_c)^{0}(4s_f)^{1}$	0.309 14	0. 281 53	0.398 20	276. 81	267. 11	264. 39
11	$(3d_c)^{3}(3d_n)^{2}(4s_c)^{1}(4s_f)^{0}$	0.300 67	0. 267 23	0.377 94	436. 31	432. 50	432.37
12	$(3d_c)^{4}(3d_n)^{2}(4s_c)^{0}(4s_f)^{0}$	0. 272 49	0. 242 33	0.342 75	440. 41	437. 19	436.94
13	$(3d_c)^{0}(3d_n)^{4}(4s_c)^{2}(4s_f)^{0}$	0. 415 41	0. 368 76	0. 521 52	65. 10	64. 39	64.39
14	$(3\mathrm{d}_{c})^{0}(3\mathrm{d}_{n})^{4}(4\mathrm{s}_{c})^{1}(4\mathrm{s}_{f})^{1}$	0. 439 69	0. 390 11	0.551 70	45. 27	44. 73	42.13
15	$(3d_c)^{1}(3d_n)^{4}(4s_c)^{0}(4s_f)^{1}$	0. 382 97	0. 340 19	0.481 12	82. 68	81.76	78.78
16	$(3d_c)^{1}(3d_n)^{4}(4s_c)^{1}(4s_f)^{0}$	0.358 80	0. 318 88	0.450 98	171.71	170. 20	170. 16
17	$(3d_c)^{2}(3d_n)^{4}(4s_c)^{0}(4s_f)^{0}$	0. 302 51	0. 269 11	0.380 65	197. 90	196. 65	196. 44

Table 2 Atom state parameters, bond parameters and characteristic properties of bcc Cr. bcc Mo. bcc W crystal

							DCC	αı,	DCC III	0, .	000 1	Cryst	·uı					
	Atom state parameter			Bond parameter							Т	Theor.		Exptl.				
	$ m d_c$	d_n	\mathbf{s}_{c}	$\mathbf{s}_{\mathbf{f}}$	R/nm	<i>r</i> ₁ / nı	n r_{j}	₂ / nm	<i>r</i> ₃ / nn	n	n_1	n_2		n 3	a/ nm	E _c / (kJmol ⁻	a/ 1) nm ^[8]	$ \begin{array}{c} E_{c}J \\ (kJ^{\bullet}) \\ \text{mol}^{-1})^{[9]} \end{array} $
bec Cr	3. 32	2. 26	0. 25	0. 17	0. 11231	0. 249	83 0.	288 4	8 0. 407	97 C	. 380 2	2 0.091	5	0.0009	0. 288 4	8 395.35	0. 288 46	395
bcc Mo	4. 23	1.48	0. 02	0. 27	0. 126 20	0. 272	540.	314 7	0 0. 445	05 C	. 461 ~	7 0.091	5	0.0006	0. 314 7	0 658.26	0. 314 66	658
bec W	5. 16	0	0. 25	0. 59	0. 130 19	0. 274	200.	316 6	20.447	77 C	. 588 :	5 0.115	5	0.000 8	0. 316 6	2 858.33	0. 316 52	859

According to the binding energies of bcc Cr, bcc Mo, bcc W and the G'(T) data of bcc Cr, hcp Cr, fcc Cr, bcc Mo, hcp Mo, fcc Mo, bcc W, hcp W, fcc W supplied in SGTE database, the binding energies of hcp Cr, fcc Cr, hcp Mo, fcc Mo, hcp W and fcc W are obtained approximately as follows:

 $E_{\rm c}({\rm hep\-Cr}) = 390.~56~{\rm kJ/mol},~E_{\rm c}({\rm fcc\-Cr}) = 387.~71~{\rm kJ/mol};~E_{\rm c}({\rm hep\-Mo}) = 646.~45~{\rm kJ/mol},~E_{\rm c}({\rm fcc\-Mo}) = 642.~80~{\rm kJ/mol},~E_{\rm c}({\rm hep\-W}) = 844.~25~{\rm kJ/mol},~E_{\rm c}({\rm fcc\-W}) = 839.~70~{\rm kJ/mol}.$

3. 2. 2 Lattice constants of hcp·Cr, fcc·Cr, hcp·Mo, fcc·Mo, hcp·W and fcc·W

As the experimental lattice constants of hcp Cr, fcc Cr, hcp Mo, fcc Mo, hcp W and fcc W can not be obtained, we use the method similar to Ref. [10], that is, supposing Cr, Mo, W metals with hcp and fcc structures have the same atom volume as those with bcc structure. Therefore, the lattice constants of hcp Cr, fcc Cr, hcp Mo, fcc Mo, hcp W and fcc W are obtained approximately as follows:

 $a(\text{hcp-Cr}) = 0.256 99 \text{ nm}, \ a(\text{fcc-Cr}) = 0.363 44 \text{ nm}; \ a(\text{hcp-Mo}) = 0.280 35 \text{ nm}, \ a(\text{fcc-Mo}) = 0.396 45 \text{ nm}; \ a(\text{hcp-W}) = 0.281 99 \text{ nm}, \ a(\text{fcc-W}) = 0.398 79 \text{ nm}.$

3. 2. 3 Electronic structures of hcp Cr, fcc Cr, hcp Mo, fcc Mo, hcp W and fcc W

Similar to determining the electronic structures of bccCr, bccMo and bccW, electronic structures of hcpCr, fccCr, hcpMo, fccMo, hcpW and fccW can be determined by the multiple properties determining state method as

$$\begin{split} \Psi_{a}(\text{hcpr}\,\text{Cr}) &= \left[\,\text{Ar}\right] (\,\,3d_{c})^{\,\,3.\,\,24} (\,\,3d_{n})^{\,\,2.\,\,40} (\,\,4s_{c})^{\,\,0.\,\,28} \bullet \\ &\quad (\,\,4s_{f})^{\,\,0.\,\,08}, \\ \Psi_{a}(\text{fcc}\,\text{Cr}) &= \left[\,\text{Ar}\right] (\,\,3d_{c})^{\,\,3.\,\,14} (\,\,3d_{n})^{\,\,2.\,\,54} (\,\,4s_{c})^{\,\,0.\,\,32}; \\ \Psi_{a}(\text{hcpr}\,\text{M}\,\text{o}) &= \left[\,\text{Kr}\right] (\,\,4d_{c})^{\,\,4.\,\,00} (\,\,4d_{n})^{\,\,1.\,\,74} \bullet \\ &\quad (\,\,5s_{c})^{\,\,0.\,\,99} (\,\,5s_{f})^{\,\,0.\,\,17}, \\ \Psi_{a}(\text{fcc}\,\text{M}\,\text{o}) &= \left[\,\text{Kr}\right] (\,\,4d_{c})^{\,\,3.\,\,79} (\,\,4d_{n})^{\,\,2.\,\,00} \bullet \\ &\quad (\,\,5s_{c})^{\,\,0.\,\,17} (\,\,5s_{f})^{\,\,0.\,\,04}; \\ \Psi_{a}(\text{hcpr}\,\text{W}) &= \left[\,\text{Xe}\right] (\,\,5d_{c})^{\,\,4.\,\,88} (\,\,5d_{n})^{\,\,0.\,\,32} (\,\,6s_{c})^{\,\,0.\,\,40} \bullet \\ &\quad (\,\,6s_{f})^{\,\,0.\,\,40}, \\ \Psi_{a}(\text{fcc}\,\text{W}) &= \left[\,\text{Xe}\right] (\,\,5d_{c})^{\,\,4.\,\,50} (\,\,5d_{n})^{\,\,0.\,\,82} (\,\,6s_{c})^{\,\,0.\,\,61} \bullet \\ &\quad (\,\,6s_{f})^{\,\,0.\,\,07}. \end{split}$$

3. 2. 4 Electronic structures of liquid metals L-Cr, L-Mo and L-W

Similar to the case of hcp Cr, fcc Cr, hcp Mo, fcc Mo, hcp W and fcc W, the binding energies of L-Cr, L-Mo and L-W are determined through SGTE database as follows:

 $E_{c}(L-Cr) = 370.66 \text{ kJ/mol}; E_{c}(L-Mo) = 616.17 \text{ kJ/mol}; E_{c}(L-W) = 806.84 \text{ kJ/mol}.$

Supposing Cr, Mo, W still have bee structure after being melted, according to their densities^[11] before and after being melted and the lattice constants of solid bec Cr, bec Mo, bec W, the lattice constants of L-Cr, L-Mo and L-W are obtained as follows:

 $a(\text{L-Cr}) = 0.30145 \text{ nm}, \ a(\text{L-Mo}) = 0.32436 \text{ nm}, \ a(\text{L-W}) = 0.32553 \text{ nm}.$ By the way, it has been confirmed by modern X-ray experiments that liquid metals really have the short distance ordering similar to crystals.

Similar to the case above, the electronic structures of L-Cr, L-Mo and L-W are obtained as: $\Psi_{a}(\text{L-Cr}) = [\text{Ar}] (3d_{c})^{3.44} (3d_{n})^{1.52} (4s_{c})^{0.20} (4s_{f})^{0.84};$ $\Psi_{a}(\text{L-Mo}) = [\text{Kr}] (4d_{c})^{4.00} (4d_{n})^{1.36} (5s_{c})^{0.08} (5s_{f})^{0.56};$ $\Psi_{a}(\text{L-W}) = [\text{Xe}] (5d_{c})^{4.69} (6s_{c})^{0.50} (6s_{f})^{0.81}.$

4 QUALITATIVE EXPLANATION OF RELA-TIONSHIP BETWEEN CRYSTALLINE STRUCTURES AND ELECTRONIC STRUC-TURES OF Cr. Mo AND W

Compared with various kinds of first-principle methods, an important feature of the description of the electronic structure in OA theory is the function division of electrons in outer shells. On the basis of this, a qualitative explanation about why the crystalline structures of metals exist can be presented.

In the full p^6 shells of Cr, Mo and W metals, electron cloud of p orbits extends along x, y and z axes. Along these directions, the shield effect of electrons to nuclear charges is the strongest, while between the axes the effect is weaker. So, there are two direction systems with weak shield effect outside p^6 shell^[5]. The first direction system is $\langle 111 \rangle$ directions along 4 body diagonal lines, and the angle between them and p_x , p_y and p_z is the biggest (54. 75°). The second is $\langle 110 \rangle$ directions along 3 pairs of diagonal lines on xy, yz and xz planes, and the corresponding angle is 45° .

The d electrons out of p⁶ shell are divided into e_g and t_{2g} , and they have different energy and symmetry: d electrons with e_g state extend along $\langle 111 \rangle$ directions (that is, the first direction system mentioned above), and those with t_{2g} extend along the $\langle 110 \rangle$ directions(that is, the second direction system mentioned above). d electrons first fill e_g orbits with low energy, and bonds along (111) directions are formed. Atoms with this state are inclined to form bcc structure crystal with low symmetry. With the increase of $d_{\rm c}$ electrons with $\it e_{\rm g}$ state , the bcc structure becomes more stable. s electrons are sphere symmetrical, their existence is favorable to forming crystalline structures with high symmetry. The coordination number of bcc structure is 8, and that of hcp and fcc structure are 6+ 6(the nearest number+ the next nearest number), 12 respectively. That is, from bcc to hep to fee, the bond directivity becomes weaker and the symmetry becomes stronger. The calculation results of electronic structures in our study can give a satisfactory explanation about why the crystalline structures of Cr, Mo, W metals exist. In the course of bcc hcp fcc, dc electrons become less gradually (Cr: 3. 32 $\stackrel{\rightarrow}{\cancel{\ }}$ 3. 24 $\stackrel{\rightarrow}{\cancel{\ }}$ 3. 14; Mo: 4. 23 $\stackrel{\rightarrow}{\cancel{\ }}$ 4. 00 $\stackrel{\rightarrow}{\cancel{\ }}$ 3. 79; W: 5. 16 $\stackrel{\rightarrow}{\cancel{\ }}$ 4. 88 $\stackrel{\rightarrow}{\cancel{\ }}$ 4. 50), while s_c electrons become more gradually (Cr: 0. 25 0. 28 0. 32; Mo: 0. 02 $0.09 \stackrel{?}{\to} 0.17$; W: $0.25 \stackrel{?}{\to} 0.40 \stackrel{?}{\to} 0.61$). This makes the symmetry of the electron cloud distribution become stronger, so the coordination number becomes

5 **QUALITATIVE EXPLANATION OF RELA-**TIONSHIP BETWEEN DIFFERENCE PHYSICAL PROPERTIES OF Cr., Mo. METALS AND THEIR ELECTRONIC STRUC-**TURES**

Based on the function division of outer-shell electrons, not only the relationship between the existence of the crystalline structures of Cr, Mo, W metals and their electronic structures, but also the relationship between the difference of physical properties of Cr, Mo, W metals and their electronic structures can be explained qualitatively. In order to compare with experimental results, only Cr, Mo and W with bcc structure are presented here.

Generally, the following two kinds of properties of metals are always taken into account: one is the mechanical properties (including elasticity, strength, tenacity, hardness and brittleness), the other is transport properties (including electrical conductivity, thermal conductivity). From the view point of condensed physics, mechanical properties are the embodiment of combination properties of metal crystals in essence. In OA theory, outershell electrons of metal atoms are divided into 4 kinds: covalent electrons n_c , near-free electrons n_f , no-bond electrons $n_{\rm n}$ and magnetic electrons $n_{\rm m}$. Covalent electrons n_c are responsible for forming the bond network to bind metal atoms. It can be considered that the more the covalent electrons are, the stronger the binding is, and the greater the strength and hardness are. Furthermore, a mechanical property parameter $M = n_c$ can be introduced to describe the binding strength of metal crystals. The M data, tensile strength, Vickers hardness, bulk modulus and melting point of Cr, Mo, W metals are listed in Table 3. It can be seen that their change trends are the same.

It must be pointed out that structure sensitive properties such as strength and hardness are connected with many factors. The organization and structure of the specimen (monocrystalline, polycrystalline, size of crystal grains), the product form of the specimen (cast product, forging part, rolled product, extruded section and P/M pro-duct), heat processing technology (quenching, annealing, annealing temperature, annealing time), the shape and size of the specimen

(rod, sheet), test conditions (strain rate, temperature) can all affect experimental results greatly. So the decentralization rate of the experimental data from various sources is great, and their comparability is bad. As a kind of structure insensitive property, the comparability of bulk modulus data is much better. Furthermore, although melting point is usually not considered as a kind of mechanical property, it can be a very good parameter to describe the binding strength of metal crystals. It can be measured precisely, and the decentralization rate of results is small. So, the melting points are also listed in Table

Table 3 Relationship between difference of mechanical properties and transport properties of Cr, Mo, W metals with bcc structure and their electronic structures

D .	M etal					
Property -	bec Cr	bcc Mo	bec W			
M	3. 57	4. 25	5. 41			
Tensile strength/MPa ^[8]	282	300	380			
Vickers hardness ^[8]	135	180	190			
Bulk modulus/ GPa ^[9]	190. 1	272. 5	323. 2			
Melting point/K ^[8]	2148	2883	3683 ± 20			
T' / nm $^{-3}$	14. 16	17. 33	37. 21			
$\sigma_{295 \text{ K}} / 10^{-5} (\Omega^{-1} \cdot \text{cm}^{-1})^{[9]}$	0.78	1.89	1. 89			
$K_{300 \text{ K}} / (\text{W} \cdot \text{cm}^{-1} \cdot \text{K}^{-1})^{[9]}$	0. 94	1.38	1. 74			

According to the solid electronic theory in condensed physics, the expressions of the electrical conductivity σ and thermal conductivity K are as follows^[9]:

$$\sigma = \frac{ne^2 T}{m}$$

$$K = \frac{\pi^2 nk_B^2 T T}{3m}$$
(5)

$$K = \frac{\pi^2 n k_B^2 T \tau}{3 m} \tag{6}$$

n is electron concentration, e is electron charge, τ is relaxation time, m is electron mass, $k_{\rm B}$ is Boltzmann constant, T is Kelvin temperature. As e, m, k_B are all constants, if we assume that the relaxation time T of Cr. Mo and W are the same, the electrical conductivity σ and thermal conductivity Kof Cr, Mo, W will be in direct proportion to the electron concentration n at the same temperature.

In OA theory, near-free electrons in outer shells of metal atoms are responsible for transporting electricity and heat. A parameter of transport property $T' = n_{\rm f}/V$ can be introduced to describe the transport property of metals, where V is the atom volume. It can be seen that the physical meaning of the transport property parameter is the near-free electrons

concentration. The transport properties parameter T', electrical conductivity Y and thermal conductivity K are also listed in Table 3. It can be seen that their change trends are also the same.

6 QUANTITATIVE CALCULATIONS OF PHYSI-CAL PROPERTIES OF Cr, Mo AND W WITH BCC STRUCTURE

Seen as a whole, one of the critical parts of OA theory is the determination of electronic structure, the other is the qualitative explanations and quantitative calculations of various physical properties based on the electronic structures obtained above. The object of researching electronic structure is to perform these kinds of explanations and calculations, and to understand the change rules of various physical properties in essence and control them effectively. On the other hand, these kinds of qualitative explanations and calculations can also be used to judge whether the electronic structures are proper or not. The potential curves, elasticity and the changes of linear thermal expansion with temperature of Cr, Mo and W with bcc structure are calculated here. It is necessary to point out that these calculations can also be performed to Cr, Mo and W with hep and fee structures and liguid state. Only calculation results of Cr, Mo and W with bcc structure are presented here, because it is convenient to compare them with experimental results.

6.1 Theoretical potential curves

According to many-atom interactions (MAI) potential function, the theoretical potential curves of Cr, Mo and W with bcc structure are shown in Fig. 1.

6. 2 Elasticity

According to the elasticity calculation formula in OA theory, bulk modulus(B), elastic modulus(Y), modulus rigidity(Y) and Poisson's ratio(Y) of bcc Cr, bcc Mo and bcc W are listed in Table 4.

6. 3 Linear thermal expansion coefficient as a function of temperature for bcc Cr, bcc Mo and bcc W

According to the expression between linear

Table 4 Elasticities of bcc Cr, bcc Mo and bcc W

	Elasticity	bcc Cr	bcc M o	
		Theor. Exptl.	Theor. Exptl.	Theor. Exptl.
	B/ GPa	184 190.1	^[9] 272 272. 5 [[]	^{9]} 316 323. 2 ^[9]
	Y/ GPa	187	276	321
	μ/ GPa	70	104	121
	Υ	0.3308	0. 3308	0. 3308

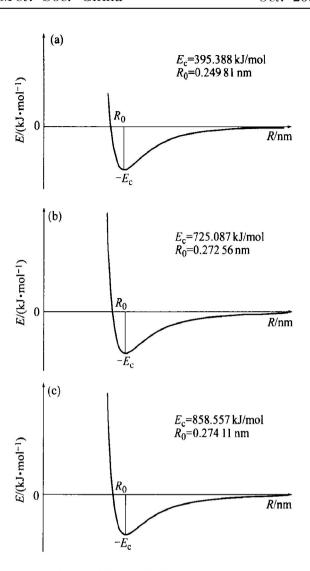


Fig. 1 Theoretical potential curves of bcc Cr(a), bcc Mo(b) and bcc W(c)

thermal expansion coefficient and temperature in OA theory, the change curves of linear thermal expansion coefficient as a function of temperature of bcc Cr, bcc Mo and bcc W are shown in Fig. 2.

7 CONCLUSIONS

- 1) When Cr, Mo and W atoms get together, there are 0.58, 0.71 and 0.16 s electrons transforming into d electrons respectively. So, the electronic structures of bcc Cr, bcc Mo and bcc W can be described as follows: $\Psi_a(bcc Cr) = [Ar] \cdot (3d_c)^{3.32} (3d_n)^{2.26} (4s_c)^{0.25} (4s_f)^{0.17}$, $\Psi_a(bcc Mo) = [Kr] (4d_c)^{4.23} (4d_n)^{1.48} (5s_c)^{0.02} (5s_f)^{0.27}$, $\Psi_a(bcc W) = [Xe] (5d_c)^{5.16} (6s_c)^{0.25} (6s_f)^{0.59}$.
- 2) Based on the electronic structures of Cr, Mo and W, the relationship between their crystalline structures and electronic structures is explained qualitatively; the relationship between the difference of the physical properties of Cr, Mo and W with bcc structure and their electronic structures is explained qualitatively; the lattice constants, binding energies,

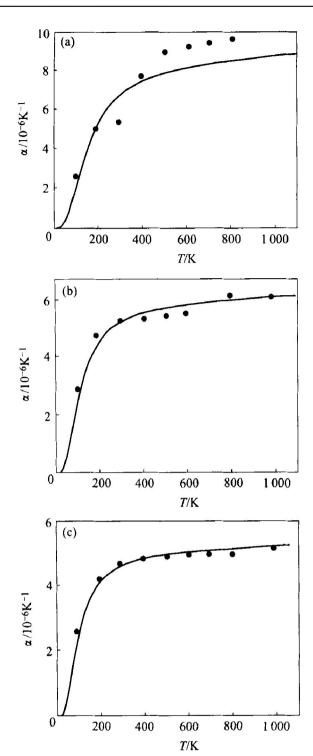


Fig. 2 Linear thermal expansion coefficient as a function of temperature for bcc Cr(a), bcc Mo(b), bcc W(c)(• —Exptl. values^[12])

potential curves, elasticity and change curves of linear thermal expansion coefficient as a function of temperature are calculated quantitatively, and the theoretical values and experimental values are in good accordance.

3) Electronic structures of hcp·Cr, fcc·Cr, L-Cr, hcp·Mo, fcc·Mo, L-Mo, hcp·W, fcc·W, L-Cr,

L-Mo and L-W are obtained as follows: $\Psi_a(\text{hcp-Cr}) = [Ar] (3d_c)^{3.24} (3d_n)^{2.40} (4s_c)^{0.28} \cdot (4s_f)^{0.08}, \quad \Psi_a(\text{fcc-Cr}) = [Ar] (3d_c)^{3.14} (3d_n)^{2.54} \cdot (4s_c)^{0.32}, \quad \Psi_a(\text{L-Cr}) = [Ar] (3d_c)^{3.44} (3d_n)^{1.52} \cdot (4s_c)^{0.32}, \quad (4s_f)^{0.84}; \quad \Psi_a(\text{hcp-Mo}) = [Kr] (4d_c)^{4.00} \cdot (4d_n)^{1.74} (5s_c)^{0.09} (5s_f)^{0.17}, \quad \Psi_a(\text{fcc-Mo}) = [Kr] (4d_c)^{3.79} (4d_n)^{2.00} (5s_c)^{0.17} (5s_f)^{0.04}, \quad \Psi_a(\text{L-Mo}) = [Kr] (4d_c)^{4.00} (4d_n)^{1.36} (5s_c)^{0.08} (5s_f)^{0.56}; \quad \Psi_a(\text{hcp-W}) = [Xe] (5d_c)^{4.88} (5d_n)^{0.32} (6s_c)^{0.40} \cdot (6s_f)^{0.40}, \quad \Psi_a(\text{fcc-W}) = [Xe] (5d_c)^{4.50} (5d_n)^{0.82} \cdot (6s_c)^{0.61} (6s_f)^{0.07}, \quad \Psi_a(\text{L-W}) = [Xe] (5d_c)^{4.69} \cdot (6s_c)^{0.50} (6s_f)^{0.81}.$

REFERENCES

- [1] XIE Yourqing. The framework of metallic materials systematic science [J]. Mater Rev, 2001, 15(4): 12 15. (in Chinese)
- [2] XIE Yourqing, MA Liurying. The theoretical lattice parameter of the valence electron structure of crystal [J]. J Cent South Inst of Mining and Metall, 1985, 16(1): 1-10. (in Chinese)
- [3] XIE Yourqing. A new potential function with many atom interactions in solid [J]. Science in China (Series A), 1992, 22(8): 880-890. (in Chinese)
- [4] XIE Yourqing. One atom self-consistency method for determining electronic structure of crystal[J]. Chinese Science Bulletin, 1992, 37(16): 1529 1532. (in Chinese)
- [5] Xie Your qing. Electronic structure and properties of pure iron[J]. Acta Metall Mater, 1994, 42 (11): 3705 -3715.
- [6] XIE Yourqing. Metallic Materials Systematic Science [M]. Changsha: Central South University of Technology Press, 1998. 32.
- [7] Dinsdale A T. SGTE data for pure elements [J]. CAL-PHAD, 1991, 15(4): 317 - 425.
- [8] ASM International Handbook Committee. Metals Handbook, 10th edition, Vol 2: Properties and Selection: Nonferrous Alloys and Special purpose Materials [M]. OH: Materials Park, 1990.
- [9] Kittel C. Introduction to Solid State Physics (6th edition)[M]. New York: John Wiley & Sons, Inc, 1986.
- [10] Barkonyi I, Elbert H, Liechtenstein A I. Electronic structure and magnetic susceptibility of the different structural modifications of Ti, Zr and Hf metals [J]. Phys Rev B, 1993, B48: 7841 – 7849.
- [11] Weast R C. CRC Handbook of Chemistry and Physics (70th edition) [M]. Florida: CRC Press Inc, 1990.
- [12] American Institute of Physics. American Institute of Physics Handbook(3rd edition) [M]. New York: McGraw-Hill Book Company, 1972.

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