

## Application of geometric midline yield criterion to analysis of three-dimensional forging

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**Abstract:** A kinematically admissible continuous velocity field was proposed for the analysis of three-dimensional forging. The linear yield criterion expressed by geometric midline of error triangle between Tresca and Twin shear stress yield loci on the  $\pi$ -plane, called GM yield criterion for short, was firstly applied to analysis of the velocity field for the forging. The analytical solution of the forging force with the effects of external zone and bulging parameter is obtained by strain rate inner product. Compression tests of pure lead are performed to compare the calculated results with the measured ones. The results show that the calculated total pressures are higher than the measured ones whilst the relative error is no more than 9.5%. It is implied that the velocity field is reasonable and the geometric midline yield criterion is available. The solution is still an upper-bound one.

**Key words:** GM yield criterion; three-dimensional forging; continuous velocity field; inner product; analytical solution

### 1 Introduction

In recent years, the researches on metal forging are focused on numerical simulations[1], including FEM [2–5] and UBEM[6]. Unfortunately, no theoretical analytical solution about three-dimensional forging taking into account of the effects of external zone has been reported in the past twenty years. So, how to substitute the non-linear Mises yield criterion with linear yield criterion to get analytical solution attracts much attention. TRESCA[7] first proposed a linear yield criterion but only took two principal stresses into account. HILL[8] also introduced a linear yield criterion to approach Mises criterion in 1950 but with a relative error by 8%. Until 1983, YU[9], HUANG and ZENG[10] proposed a linear criterion, called Twin shear stress yield criterion, but its application to analysis of metal forging [11–12] usually shows a greater value than that of Mises yield criterion. Based on the error triangles consisting of Tresca and Twin shear stress yield loci on the  $\pi$ -plane, the GM yield criterion was proposed by ZHAO et al[13–15].

In this work, a new velocity field was put forward to the three-dimensional forging and the GM yield criterion was firstly applied to analysis of the forging. The solution was compared with the tested result.

### 2 Velocity field

Three-dimensional forging with external zone between two parallel indenters is shown as Fig.1.

Because of symmetry, only one eighth of the deformation zone is taken into account and it is assumed

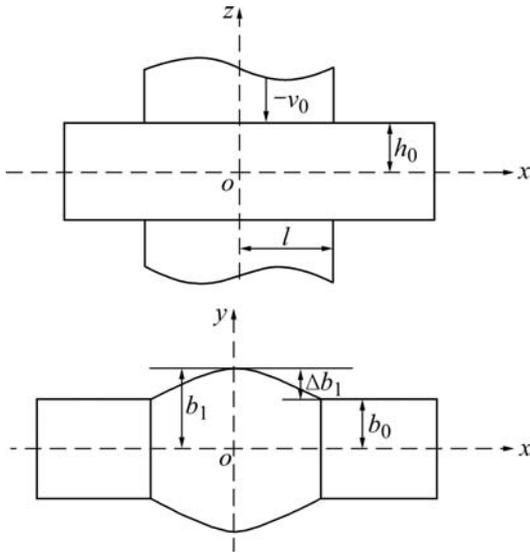
$$\Delta b_1/b_0 = a\Delta h/h_0 \quad (1)$$

Dividing the both sides of Eqn.(1) by time increment  $\Delta t$  yields

$$\frac{v_{b_1}}{b_0} = a \frac{v_0}{h_0} \quad (2)$$

Let the velocity component  $v_z$  vary linearly with  $z$  coordinate, then it follows

$$v_z = -\frac{v_0}{h_0} z \quad (3)$$



**Fig.1** Deformation zone of 3-dimensional forging

Assuming that in Fig.1 the free-edge surface is parabolic, the width  $b$  is

$$b = b_0 + \Delta b_1 \left(1 - \frac{x^2}{l^2}\right) \quad (4)$$

where  $\Delta b_1 = b_1 - b_0$  is measured at the maximum width point. Dividing Eqn.(4) by  $\Delta t$  yields

$$v_b = v_{b_1} \left(1 - \frac{x^2}{l^2}\right) \quad (5)$$

Let  $v_y$  vary linearly with  $y$  coordinate, which leads to

$$v_y = \frac{v_b}{b_0} y = \frac{v_{b_1}}{b_0} \left(1 - \frac{x^2}{l^2}\right) y = a \frac{v_0}{h_0} \left(1 - \frac{x^2}{l^2}\right) y \quad (6)$$

From the Cauchy equation it follows

$$\begin{cases} \dot{\epsilon}_y = \frac{\partial v_y}{\partial y} = a \frac{v_0}{h_0} \left(1 - \frac{x^2}{l^2}\right) \\ \dot{\epsilon}_z = \frac{\partial v_z}{\partial z} = -\frac{v_0}{h_0} \\ \dot{\epsilon}_x = -(\dot{\epsilon}_y + \dot{\epsilon}_z) = \frac{v_0}{h_0} \left[1 - a \left(1 - \frac{x^2}{l^2}\right)\right] \\ \dot{\epsilon}_{xy} = \frac{1}{2} \left(\frac{\partial v_y}{\partial x}\right) = -a \frac{v_0 \cdot xy}{h_0 l^2} \\ \dot{\epsilon}_{xz} = \dot{\epsilon}_{yz} = 0 \end{cases} \quad (7)$$

From  $v_x = \int \dot{\epsilon}_x dx + \psi(y, z)$ , and letting  $x=0$ ,  $v_x=0$  yields

$$\begin{cases} v_x = \frac{v_0}{h_0} x \left[1 - a \left(1 + \frac{x^2}{3l^2}\right)\right] \\ v_y = a \frac{v_0}{h_0} \left(1 - \frac{x^2}{l^2}\right) y \\ v_z = -\frac{v_0}{h_0} z \end{cases} \quad (8)$$

Note that in Eqn.(8),  $y=0$ ,  $v_y=0$ ;  $y=b_0$ ,  $v_y=v_b$ ;  $z=0$ ,  $v_z=0$ ;  $z=h_0$ ;  $v_z=-v_0$ ; and in Eqn.(7),  $\dot{\epsilon}_x + \dot{\epsilon}_y + \dot{\epsilon}_z = 0$ . So, Eqns.(7) and (8) satisfy kinematically admissible conditions.

With integral mean value theorem and the volume constancy, it follows

$$\begin{cases} \bar{v}_y = \frac{1}{b_0 l} \int_0^{b_0} \int_0^l v_y dx dy = \frac{av_0 b_0}{3h_0} \\ \bar{v}_x = \frac{lv_0}{2h_0} \left(1 - \frac{5a}{6}\right) \\ \bar{\dot{\epsilon}}_x = \frac{1}{lb_0} \int_0^{b_0} \int_0^l \dot{\epsilon}_x dx dy = \frac{v_0}{h_0} \left(1 - \frac{2a}{3}\right) \\ \bar{\dot{\epsilon}}_y = \frac{2av_0}{3h_0} \\ \bar{\dot{\epsilon}}_{xy} = \frac{av_0 b_0}{4lh_0} \end{cases} \quad (9)$$

Values of Eqn.(9) are used to calculate the ratios of components of strain rate and velocity field.

The characteristic equation of strain rate tensor has a non-vanishing solution only if the following determinant vanishes, that is

$$\begin{vmatrix} \dot{\epsilon}_x - \dot{\epsilon} & \dot{\epsilon}_{xy} & 0 \\ \dot{\epsilon}_{yx} & \dot{\epsilon}_y - \dot{\epsilon} & 0 \\ 0 & 0 & \dot{\epsilon}_z - \dot{\epsilon} \end{vmatrix} = 0$$

Therefore, the principal strain rate field is

$$\begin{cases} \dot{\epsilon}_{1,2} = \pm \sqrt{\dot{\epsilon}_{xy}^2 - \dot{\epsilon}_x \dot{\epsilon}_y + \frac{1}{4} \dot{\epsilon}_z^2 - \frac{\dot{\epsilon}_z}{2}} = \dot{\epsilon}_{\max,2} \\ \dot{\epsilon}_3 = \dot{\epsilon}_z = \dot{\epsilon}_{\min} \end{cases} \quad (10)$$

### 3 GM yield criterion

So-called GM yield criterion is a short of geometric midline yield criterion, whose yield locus on  $\pi$ -plane is the geometric midline  $B'E$  of error triangle  $B'BF$  between Tresca locus  $B'F$  and Twin shear stress locus  $B'B$ . It intersects with Mises circle as shown in Fig.2. The detail of the criterion can be seen in Refs.[12–14]. Its plastic work rate done per unit volume is as follows[12]:

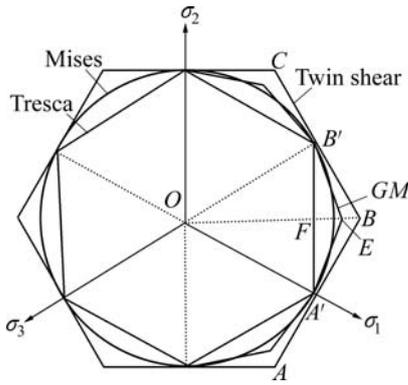


Fig.2 GM loci on  $\pi$ -plane

$$D(\dot{\epsilon}_{ij}) = \frac{7}{12} \sigma_s (\dot{\epsilon}_{\max} - \dot{\epsilon}_{\min}) \tag{11}$$

**4 Total power and pressure**

**4.1 Plastic work rate( $N_d$ )**

Substituting Eqn.(10) into Eqn.(11) and the following equation, and from Eqn.(7), it yields

$$\begin{aligned} N_d &= \int_V D(\dot{\epsilon}_{ij}) dV \\ &= \frac{7\sigma_s}{12} \int_V (\dot{\epsilon}_{\max} - \dot{\epsilon}_{\min}) dV \\ &= \frac{7\sigma_s}{12} \int_V \left( \sqrt{\dot{\epsilon}_{xy}^2 - \dot{\epsilon}_x \dot{\epsilon}_y + \frac{1}{4} \dot{\epsilon}_z^2} - \frac{3}{2} \dot{\epsilon}_z \right) dV \\ &= \frac{7h_0\sigma_s}{12} \left\{ \int_0^{b_0} \int_0^l \frac{|\dot{\epsilon}_{xy}| dx dy}{\sqrt{1 - \frac{\dot{\epsilon}_x}{\dot{\epsilon}_{xy}} \frac{\dot{\epsilon}_y}{\dot{\epsilon}_{xy}} + \frac{1}{4} \left( \frac{\dot{\epsilon}_z}{\dot{\epsilon}_{xy}} \right)^2}} - \right. \\ &\quad \left. \int_0^{b_0} \int_0^l \frac{\dot{\epsilon}_x dx dy}{\sqrt{\left( \frac{\dot{\epsilon}_{xy}}{\dot{\epsilon}_y} \right) - \frac{\dot{\epsilon}_x}{\dot{\epsilon}_y} + \left( \frac{\dot{\epsilon}_z}{2\dot{\epsilon}_y} \right)^2}} + \right. \\ &\quad \left. \int_0^{b_0} \int_0^l \frac{|\dot{\epsilon}_z| dx dy}{4 \sqrt{\left( \frac{\dot{\epsilon}_{xy}}{\dot{\epsilon}_z} \right)^2 - \frac{\dot{\epsilon}_x}{\dot{\epsilon}_z} \frac{\dot{\epsilon}_y}{\dot{\epsilon}_z} + \frac{1}{4}}} \right\} + \\ &= \frac{7h_0\sigma_s}{8} \int_0^{b_0} \int_0^l \frac{v_0}{h_0} dx dy \\ &= \frac{7h_0\sigma_s}{12} (I_1 - I_2 + I_3) + \frac{7}{8} \sigma_s v_0 b_0 l \tag{12} \end{aligned}$$

where  $I_1, I_2$  and  $I_3$  are the termwise integrations of the

inner-product of strain rate vector[15–17]. Substituting the values of Eqn.(9) into the denominator of the following fraction and integrating leads to

$$\begin{cases} I_1 = \int_0^{b_0} \int_0^l \frac{|\dot{\epsilon}_{xy}| dx dy}{\sqrt{1 - \frac{\bar{\epsilon}_x}{\dot{\epsilon}_{xy}} \frac{\bar{\epsilon}_y}{\dot{\epsilon}_{xy}} + \frac{1}{4} \left( \frac{\bar{\epsilon}_z}{\dot{\epsilon}_{xy}} \right)^2}} \\ = \frac{3a^2 b_0^3 v_0}{4h_0 \sqrt{(3ab_0)^2 - 32a^2(3-2a) + (6l)^2}} \\ I_2 = \frac{8l^2 a v_0 (3-2a) b_0}{3h_0 \sqrt{(3ab_0)^2 - 32a^2(3-2a) + (6l)^2}} \\ I_3 = \frac{3l^2 v_0 b_0}{h_0 \sqrt{(3ab_0)^2 - 32a^2(3-2a) + (6l)^2}} \end{cases}$$

Substituting  $I_1, I_2$  and  $I_3$  into Eqn.(12) yields

$$\begin{aligned} N_d &= \frac{7}{8} \sigma_s v_0 b_0 l + \frac{7\sigma_s v_0 b_0 l}{144} \\ &\quad \sqrt{(3a)^2 \frac{b_0^2}{l^2} - 32a(3-2a) + 36} \tag{13} \end{aligned}$$

**4.2 Friction power( $N_f$ )**

At contact interface, let

$$\tau_f = \frac{m\sigma_s}{\sqrt{3}} \tag{14}$$

and

$$|\Delta v_f| = v_f = \sqrt{v_x^2 + v_y^2} \tag{15}$$

$$\begin{aligned} N_f &= \int_s \tau_f |\Delta v_f| ds = \frac{m\sigma_s}{\sqrt{3}} \int_0^{b_0} \int_0^l \sqrt{v_x^2 + v_y^2} dx dy \\ &= \frac{m\sigma_s}{\sqrt{3}} \int_0^{b_0} \int_0^l \left[ 1 + \left( \frac{v_y}{v_x} \right)^2 \right]^{1/2} dx dy \tag{16} \end{aligned}$$

Substituting  $\bar{v}_x/\bar{v}_y$  of Eqn.(9) for  $v_x/v_y$  into Eqn.(16), it follows

$$N_f = \frac{m\sigma_s l b_0 v_0}{12\sqrt{3}h_0} \sqrt{l^2(6-5a)^2 + (4ab_0)^2} \tag{17}$$

**4.3 Shear power ( $N_s$ )**

At the interface between deforming and external zones, the velocity discontinuity and shear power are

$$\begin{cases} \Delta v_t = \bar{v}_z = \frac{1}{h} \int_0^{h v_0} z dz = \frac{v_0}{2} \\ N_s = k b_0 h_0 \Delta v_t = \frac{\sigma_s b_0 h_0 v_0}{2\sqrt{3}} \end{cases} \tag{18}$$

#### 4.4 Total power and pressure

Let applied power equal to the upper bound total power, we can have

$$\bar{p}v_0b_0l = N_d + N_f + N_s \quad (19)$$

Substituting Eqns.(13), (17) and (18) into the equation above and rearranging yields

$$n_\sigma = \frac{\bar{p}}{\sigma_s} = \frac{7}{8} + \frac{m}{12\sqrt{3}} \sqrt{(6-5a)^2 \frac{l^2}{h_0^2} + \left(\frac{4ab_0}{h_0}\right)^2} + \frac{7}{48} \sqrt{\left(\frac{ab_0}{l}\right)^2 - \frac{32}{3}a\left(1-\frac{2}{3}\right)a + 4 + \frac{h_0}{2\sqrt{3}l}} \quad (20)$$

Noting symmetry of the deforming zone, the total pressure becomes

$$P = \bar{p}S = 4n_\sigma\sigma_s\bar{b}l \quad (21)$$

Letting  $dn_\sigma/da=0$  in Eqn.(20) gives

$$m = \frac{\frac{28}{\sqrt{3}} - \frac{7\sqrt{3}a}{4} \frac{b_0^2}{l^2} - \frac{112\sqrt{3}a}{9}}{16a \frac{b_0^2}{h_0^2} - (30-25a) \frac{l^2}{h_0^2}} \sqrt{\frac{(6-5a)^2 \frac{l^2}{h_0^2} + 16a^2 \frac{b_0^2}{h_0^2}}{a^2 \frac{b_0^2}{l^2} - \frac{32a}{3} + \frac{64a^2}{9} + 4}} \quad (22)$$

Eqn.(22) is the relationship among  $a$ ,  $m$ ,  $l/h_0$  and  $b_0/h_0$ , which can be used for optimization of  $a$ .

The value of the friction factor  $m$  can also be calculated by the following Tarnovskii equation:

$$m = f + \frac{1}{8} \frac{l}{h_0} (1-f)\sqrt{f} \quad (23)$$

where  $f$  is coulomb friction coefficient.

In this work the authors suggest the following method to measure value of  $a$ . Substituting  $x=0$  and  $y=b_0$  into the second formula in Eqn.(8) and multiplying  $\Delta t$  yields

$$v_y \Big|_{x=0, y=b_0} \Delta t = \frac{av_0b_0\Delta t}{h}, \quad \frac{\Delta b_1}{b_0} = \frac{a\Delta h}{h}, \quad a = \frac{\Delta b_1}{b_0\varepsilon} \quad (24)$$

where  $\Delta b_1=b_1-b_0$  is the measured spread at the maximum width point, and  $\varepsilon$  is the reduction in the pass.

#### 5 Verification by press test

The press tests were performed with 200 kN universal material testing machine in the State Key Laboratory of Rolling and Automation, Northeastern University. Four groups of pure lead sample were compressed with different indenters and reductions. The indenter speeds were from 15 to 30 mm/min. The sample size and tested data are listed in Table 1.  $P_m$  is the measured total pressure.

Calculated results according to Eqn.(20) with measured values of  $a$  by Eqn.(24) are listed in Table 2. Taking the No.2 as an example, the detailed procedure is as follows: from Table 1 and Eqn.(24),  $\varepsilon=(9.85-8.745)/9.85=11.2\%$ ,  $l/h_0=1.52$ ,  $a=0.51$ ,  $f=0.23$  (for quenched heads). By Eqn.(23),  $m=0.3$ . Substituting all data of No.2 into Eqns.(20) and (21) yields  $n_\sigma=1.296$ .

$P=32.31$  kN, and the relative error with  $P_m$  is  $\Delta=(32.31-30)/30=7.7\%$ . In above calculation  $\sigma_s=20.26$  MPa, is checked out by  $\dot{\varepsilon}=0.112/t=0.025$  s<sup>-1</sup>,  $\varepsilon=11.2\%$ . Calculations of the other specimen are the same.

It can be seen from Table 2 that the total pressure  $P$ , calculated by Eqn.(20), are greater than measured ones  $P_m$ . Both relative errors get to 7.5%–18.3%.

It can be seen from Table 3 that the optimized total pressures  $P$ , by golden mean according to Eqns.(20) and (22), are lower than those in Table 2. Whilst the relative

**Table 1** Sample size and tested data

Test No.	$h_0$ /mm	$b_0$ /mm	$l$ /mm	$h_1$ /mm	$b_1$ /mm	$P_m$ /kN
1	10.165	10.165	15.0	9.250	10.85	14.1
2	9.850	19.930	15.0	8.745	21.08	30.0
3	5.035	15.025	7.5	4.325	15.59	11.9
4	5.060	19.990	7.5	5.540	20.38	14.5

**Table 2** Calculated results by measured  $a$  ( $m=0.3$ )

Test No.	$b_0/h_0$	$a$	$n_\sigma$	$P$ /kN	$P_m$ /kN	$\Delta$ /%
1	1	0.75	1.210	15.44	14.1	9.5
2	2	0.51	1.296	32.31	30.0	7.7
3	3	0.27	1.380	12.80	11.9	7.5
4	4	0.19	1.410	17.15	14.5	18.3

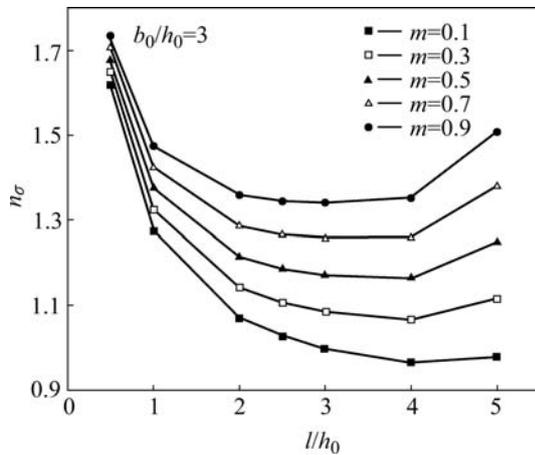
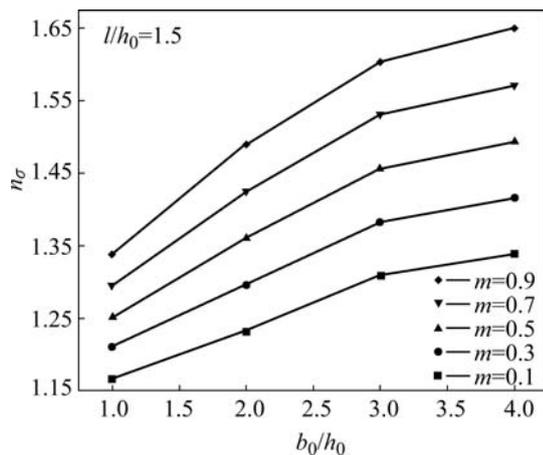
**Table 3** Optimized results by golden mean ( $m=0.3$ )

Test No.	$b_0/h_0$	$a$	$n_\sigma$	$P/\text{kN}$	$P_m/\text{kN}$	$\Delta/\%$
1	1	0.72	1.120	14.44	14.10	2.4
2	2	0.60	1.201	30.21	30.0	0.4
3	3	0.45	1.270	11.90	11.9	0
4	4	0.33	1.310	15.88	14.5	9.5

errors of  $P$  and  $P_m$  are reduced to 0–9.5%.

Fig.3 shows the calculation curve corresponding to  $b_0/h_0=3$ . For a given  $l/h_0$ ,  $n_\sigma$  increases with increase of  $m$ , but for definite  $m$  and  $l/h_0$  from 1 to 5, it always has a minimal value.

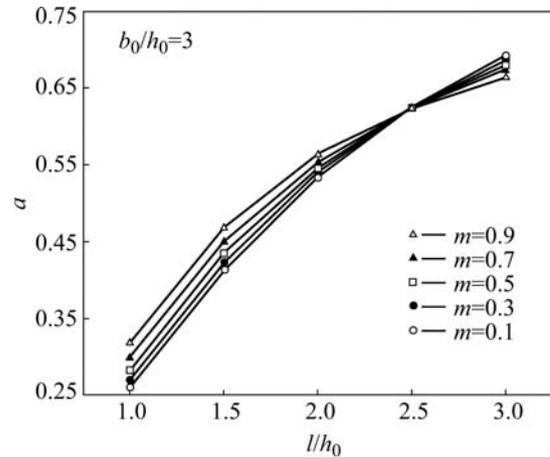
Fig.4 shows that the value of  $n_\sigma$  increases with increase of  $m$  and  $b_0/h_0$  for a given  $l/h_0$ .

**Fig.3** Dependence of  $n_\sigma$  on  $m$  and  $l/h_0$  with  $b_0/h_0=3$ **Fig.4** Dependence of  $n_\sigma$  on  $b_0/h_0$  and  $m$ 

The dependence of  $a$  on  $l/h_0$ ,  $b_0/h_0$  and  $m$  is shown in Fig.5. It can be seen that for a given  $m$ ,  $a$  increases with increase of  $l/h_0$ .

## 6 Conclusions

1) The velocity and strain rate fields proposed satisfy kinematically admissible condition of three-

**Fig.5** Dependence of  $a$  on  $l/h_0$  and  $m$ 

dimensional forging.

2) With the velocity and strain rate fields, the GM linear yield criterion is first applied to analysis of three-dimensional forging and an analytical solution of  $n_\sigma$  is obtained.

3) With press test, the calculated total pressures are higher by 7.5%–18.3%, but optimized total pressures are higher only by 0–9.5%, compared with the measured results. It is still an upper-bound solution.

4) Formula of measuring bulge parameter  $a$  is presented, and the measured values of  $a$  are lower than optimized ones.

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